

Wavelet Techniques for Solving Proportional Delay Differential Equations

Thesis submitted for the award of the Degree of Doctor of Philosophy

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by

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Abbreviations

BVPs	Boundary Value Problems
DDEs	Delay Differential Equations
FDEs	\mathbf{F} ractional \mathbf{D} ifferential \mathbf{E} quations
\mathbf{FT}	Fourier Transform
HWSM	Haar Wavelet Series Method
IVPs	Initial Value Problems
IRKHSM	Iterative Reproducing Kernel Hilbert Space Method
MAEs	Maximum Absolute Errors
RMS	Root Mean Square
STFT	Short Time Fourier Transform
MRA	\mathbf{M} ultiresolution \mathbf{A} nalysis
PDRDEs	$ {\bf P} {\rm roportional} \ {\bf D} {\rm elay} \ {\bf R} {\rm iccati} \ {\bf D} {\rm ifferential} \ {\bf E} {\rm quations} $
R-L	Riemann Liouville
RDEs	\mathbf{R} iccati \mathbf{D} ifferential \mathbf{E} quations
VIM	\mathbf{V} aritional Iteration \mathbf{M} ethod

Chapter 1

Introduction

1.1 Historical background of Wavelet

Wavelet theory appeared with independent discoveries from Morlet in the context of signal processing of seismic data and Calderon [21] in mathematical analysis. The first proper formulation of wavelet started to appear more frequently in the 1980s; however similar ideas can be traced back to the work of Hungarian mathematician Alfred Haar in 1910. The first wavelet transformation was born in 1910 and named the Haar function. It was defined as a short negative pulse followed by a short positive pulse. A few decades later, in 1946, D. Gabor invented a wavelet, whose method is very similar to the Fourier transformation, where the principle is to apply a window defined by a Gaussian function. Until the 1980s, it was not a question of talking about wavelet transformation for the functions of A. Haar and by D. Gabor [105, 31]. But it was in 1982 that the wavelet method was for first time introduced by J. Morlet, and two years later, Grossman and Morlet [43] introduced the concept of decomposition of any arbitrary square-integrable function into square integrable wavelets of constant shape and laid the foundation of wavelet methods in mathematics. Another significant development in the context of wavelet theory is multi-resolution analysis (MRA); this remarkable technique is proposed by Mallat and Meyer [70] and explains the general formalism for the construction of an orthogonal wavelet basis. In 1987, I. Daubechies created orthogonal wavelets, identified under the name of Daubechies wavelets [28]. They are characterized by their number of zero moments. Moreover, different fields can easily use this family of wavelets. The wavelet method is a potent tool in its original domain (seismology); this technique's performance has aroused the interest of different researchers to extrapolate the theory to other fields such as image compression, medical imaging, video telephony, and finite element modelling. This work sheds light on the wavelet method in numerical analysis. In this context, this thesis is mainly devoted to solve a special class of delay differential equations known as the pantograph equation. Haar wavelet series and delayed Haar wavelet series have been used to solve linear and nonlinear initial value problems of Integer and fractional order, Boundary value problems, and system of differential equations. But before looking specifically at this technique, it is essential to understand the wavelet method with its particularities, projection spaces and definition, and existing wavelet families.

The term "wavelet" is translated from the French word "Ondelette", meaning "small wave". Wavelet functions are known to serve as the bases for square integrable functions and hence are used to decompose square-integrable functions at multiple scales. Therefore, wavelets make it possible to analyze a signal in the time-frequency domain. This precise analysis makes it possible to extract distinctive elements to attenuate parasitic noises and thus represent a signal broken down into several spaces. Whereas Fourier transform decomposes a function(signal) in terms of trigonometric series. The standard Fourier transform defines the representation of a function(signal) as the sum or integral of the periodic sine and cosine functions. Fourier transform of a function gives us the frequency spectrum of that function which shows what frequencies exist in the function(signal). The Fourier transform of a function is also called the Fourier spectrum of that function [4, 26, 28, 31, 105].

1.2 Short Time Fourier Transform

Fourier transform is a useful tool to determine the frequency characteristics of a function(signal), but it is not suitable if the function(signal) has a time-varying frequency. It has been known for some time, nevertheless, that the global Fourier transform is not appropriate for computing the spectrum information of a function (signal), as it requires all past and present information of the signal to determine its spectral density at a single frequency. Window Fourier Transform (WFT) or short-time Fourier Transform (STFT) have both been developed to address this issue. To localise the time-frequency of a non-stationary function(signal), a window function is defined. The window function's width must match the portion of the signal that can be considered stationary. Segments of the function's spectral information can be ascertained by moving the window function along the time axis. Again, the problem arises due to the fixed width of the window function, and hence wavelet transform comes into play. Before going into the technicalities of wavelets, we will provide some mathematical definitions related to STFT [4, 105].

Definition 1.1. [105, 4] A non trivial function $g \in L^2(\mathbb{R})$ is called a **window** function if $t * g(t) \in L^2(\mathbb{R})$. That is the function decays to zero rapidly.

The two most important parameters for a window functions are its radius and center.

Definition 1.2. [4] The center t^* and root mean square (RMS) radius, Δg of window function g are defined by

$$t^* = \frac{1}{\|g\|_2^2} \int_{-\infty}^{\infty} t|g(t)|^2 dt, \qquad (1.2.1)$$

and

$$\Delta g = \frac{1}{\|g\|_2^2} \left(\int_{-\infty}^{\infty} (t - t^*) |g(t)|^2 dt \right)^{1/2}, \qquad (1.2.2)$$

respectively. The width of the window function g will be $2\Delta g$.

Similarly, we can have as frequency window $\hat{g}(\omega)$ with center ω^* and RMS radius $\Delta \hat{g}$ defined analogous to (1.2.1), (1.2.2). D. Gabor was the first who realised the drawback of Fourier analysis and brought up the idea of window function and short-time Fourier transform (STFT). He combines the Fourier transform and Gaussian distribution function to produce a spectrogram that plots frequency against time. Gabor transform, also known as STFT, is nothing but a Fourier transform of a function weighted by the Gaussian window, which is sliding across in time. Unlike Fourier transform, Gabor transform tells about the location of a certain frequency component.

Definition 1.3. [4] The Gabor transform of a function f in the time-frequency domain (s, ω) is denoted by $G(f)(s, \omega)$, and is defined as

$$G(f)(s,\omega) = \hat{f}_g(s,\omega) = \int_{-\infty}^{\infty} f(t)\overline{g(t-s)}dt e^{-i\omega t}$$
(1.2.3)

where $g(t) = e^{-(t-s)^2/a^2}$ is Gaussian distribution function. The sliding process of g is controlled by the parameter s (s is the center of g), and the parameter a determine the spread of the short-time window for the Fourier transform[18]. If we look closely, the Gabor transform is effectively a convolution of the function (signal) $f(t)e^{-i\omega t}$ with the function g(t).

Definition 1.4. [4] If a function $f \in L^2(\mathbb{R})$, then inversion formula

$$f(t) = G^{-1}(\widehat{f}_g(s,\omega)) = \frac{1}{2\pi \|g\|_2^2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \widehat{f}_g(s,\omega) \overline{g(t-s)} e^{i\omega t} d\omega dt.$$

Properties of Gabor Transform: Let $f, g, h \in L^2(\mathbb{R})$ and c, d be any two arbitrary constants. Then

a. Linearity:

$$G_g[cf + dh](s, \omega) = cG_gf(s, \omega) + dG_gh(s, \omega).$$

b. Time shift:

$$G_g[Tdf](s,\omega) = G_g[f(t-d)](s,\omega),$$

$$= e^{-id\omega}G_gf(s-d,\omega).$$

c. Exponential Modulation:

$$G_g[M_d f](s, \omega) = G_g[e^{id\omega} f(t)](s, \omega),$$

= $G_g f(t)(s, \omega - d),$
= $T_d G_f(s, \omega).$

d. Conjugation:

$$G_g f(s,\omega) = \overline{G_g f(s,-\omega)}.$$

1.3 Uncertainty Principle

There is a fundamental uncertainty principle in time-frequency analysis that limits the capacity to achieve high resolution in both the time and frequency domains at the same time. The degree of detail we can see in each area is intuitively recognized as resolution. The resolution depends upon the width of the window function. A small window produces a good resolution in time but gives poor frequency resolution, while a large window shows opposite behaviours; that is, if we want high resolution of the frequency content of function(signal), we have to compromise with the time at which the frequency component occurs. Hence, there is always a trade-off between simultaneous time and frequency localisation.

Definition 1.5. Let $f \in L^2(\mathbb{R})$, the dispersion of f about the point $d \in \mathbb{R}$ is the quantity

$$\Delta_d f = \frac{\int_{-\infty}^{\infty} (t-d)^2 \|f(t)\|^2 dt}{\int_{-\infty}^{\infty} \|f(t)\|^2 dt}$$
(1.3.1)

The dispersion about a point "d" is the measure of deviation or spread of its graph from t = d. This dispersion will be small if the graph of f is concentrated near t = d and is spread out away from t = d. In frequency domain,

$$\Delta_r \hat{f} = \frac{\int_{-\infty}^{\infty} (t-d)^2 \|f(t)\|^2 dt}{\int_{-\infty}^{\infty} \|f(t)\|^2 dt}.$$
(1.3.2)

Theorem 1.6. [4] (Uncertainty principle) Suppose f is a function in $L^2(\mathbb{R})$ which vanishes at $\pm \infty$. Then

$$\Delta_d f \cdot \Delta_r \hat{f} \ge \frac{1}{4}$$

for all points $d, r \in \mathbb{R}$.

The statement implies that $\Delta_d f$, $\Delta_r \hat{f}$ cannot simultaneously be small. In other words, when the time-frequency cell is narrow in time it is wider in frequency and vice-versa. In case of Gaussian function $f(t) = \frac{e^{\frac{-t^2}{2\sigma^2}}}{\sqrt{2\pi\sigma}}$ equality is achieved.

1.4 Axioms of Multiresolution Analysis and Wavelet Transform

[105] Multiresolution analysis (MRA) is a new and remarkable concept of wavelet theory. It provides a general framework using which one can construct their own families of orthogonal wavelet basis. Mallat's idea of MRA is to approximate a function at different resolutions by projecting them into different spaces.

Definition 1.7. MRA of $L^2(\mathbb{R})$ consists of a sequence $\{V_j\}_{j\in\mathbb{Z}}$ of nested closed subspaces of $L^2(\mathbb{R})$ satisfying,

$$V_j \subset V_{j+1} \ \forall j \in \mathbb{Z}.$$
$$\overline{\bigcup_{j=0}^{\infty} V_j} = L^2(\mathbb{R}), \cap_{j \in \mathbb{Z}} V_j = \{0\}.$$

Furthermore, for a function $f \in L^2(\mathbb{R})$, it is required that

$$f(t) \in V_0 \iff f(2^j t) \in V_j \forall j \in \mathbb{N}$$
, (invariance to dilation)

 $\{\phi(t-k): k \in \mathbb{Z}\}$ is an orthonoraml basis for V_0 (invariance to translation),

where $\phi(t) \in V_0$ is called scaling function.

This definition has the consequence that at each degree of resolution j, the family of functions $\{\phi_{j,k} : t \longrightarrow 2^j \phi(2^j t - k)\}_{k \in \mathbb{Z}}$ forms an orthonormal basis of the space V_j with respect to L^2 -norm. As ϕ belongs to V_0 , which is included in V_1 , it follows that ϕ can be expressed as a linear combination of $\{\phi_{1,k}\}_{k \in \mathbb{Z}}$. In other words, there exists a sequence of reals $(h_k)_{k \in \mathbb{Z}}$ such that

$$\forall t \in \mathbb{R}, \phi(t) = \sum_{k \in Z} h_k \phi(2t - k).$$

The principle of this relation, called two-scale relation, allows us to develop fast decomposition or reconstruction algorithms in the context of a multiresolution analysis. It is, therefore, interesting to be able to refine the knowledge of a function by increasing the level of resolution without recalculating all of the associated coefficients. In V_{j+1} there exist a detail space W_j serves as the orthogonal complement of V_j in V_{j+1} . i.e.

$$V_{j+1} = V_j \oplus W_j. \tag{1.4.1}$$

The space W_j includes all the functions in V_{j+1} that are orthogonal to all those in V_j under some chosen norm. The set of functions which forms the basis for W_j are called wavelets. An immediate consequence of definition 1.7 and equation (1.4.1) for any $j_0 \in \mathbb{Z}$, the space $L^2(\mathbb{R})$ verifies:

$$L^{2}(\mathbb{R}) = \overline{V_{j_{0}} \oplus \bigoplus_{j=j_{0}}^{\infty} W_{j}}$$
(1.4.2)

There exists a function ψ such that $\{t \longrightarrow \psi(t-k)\}_{k \in \mathbb{Z}}$ is an orthonormal basis of W_0 . The function ψ is called a wavelet. Again, for different level of resolution $j \in \mathbb{Z}$, the family $\{\psi_{j,k} : t \longrightarrow 2^{j/2}\psi(2^jt-k)\}_{k \in \mathbb{Z}}$ forms an orthonormal basis of the W_j space. Moreover, from the fact that space W_0 is included in V_1 , the following two-scale relation can be established:

$$\forall t \in \mathbb{R}, \psi(t) = \sum_{k \in \mathbb{Z}} g_k \phi(2t - k)$$
(1.4.3)

where $(g_k)_{k\in\mathbb{Z}}$ is sequence of real numbers. A simple example of multiresolution analysis is that of Haar, generated by the scale function $\phi = \mathbb{I}[0, 1)$ and the wavelet $\psi = \mathbb{I}_{[1/2,1)} - \mathbb{I}_{[0,1/2)}$, where \mathbb{I} is characteristics or indicator function. Using the decomposition of the $L^2(\mathbb{R})$ space given in MRA definition (1.7) one can deduce that any function f belonging to $L^2(\mathbb{R})$ is written as:

$$f = \sum_{k \in \mathbb{Z}} \alpha_{j_0,k} \phi_{j_0,k} + \sum_{j=j_0}^{\infty} \sum_{k \in \mathbb{Z}} \beta_{j,k} \psi_{j,k},$$

where $\alpha_{j_0,k} = \int f \phi_{j_0,k}$ and $\beta_{j_0,k} = \int f \psi_{j,k}$.

In signal analysis, wavelet decomposition is a frequently used technique. Its key benefit is the ability to track the temporal evolution of a signal's frequency content. The analysis of non-stationary signals can therefore be done using it instead of the Fourier transform. In mathematics, a wavelet ψ is a summable square function of the Hilbert space $L^2(\mathbb{R})$, with an oscillating evolution most of the time and a zero average. This function is often chosen as a multiscale analysis, and reconstruction tool and many problems have been solved using the wavelet method. This function ψ is called wavelet if it satisfies the following admissibility condition in the frequency domain:

$$\int_{\mathbb{R}^+} \frac{|\hat{\psi}(\omega)|^2}{|\omega|} d\omega = \int_{\mathbb{R}^-} \frac{|\hat{\psi}(\omega)|^2}{|\omega|} d\omega < +\infty.$$

Note that $\hat{\psi}$ denotes the Fourier transform of ψ . This leads to the condition that the wavelet has zero integral:

$$\int_R \psi(t)dt = 0$$

This condition invoked that the ψ has its zeroth moment vanishing. Also, ψ admits k zero moment if :

$$\int_R t^k \psi(t) dt = 0, \text{ where } k = 0, 1, \dots k.$$

Definition 1.8. [105] A double index family of wavelets generated by translation and dilation of ψ is:

$$\psi_{a,b}(t) = \frac{1}{\sqrt{|a|}} \psi\left(\frac{t-b}{b}\right), \ a, b \in \mathbb{R}, a \neq 0.$$
(1.4.4)

Definition 1.9. [105] The **continuous wavelet transform** of a function(signal) *f* can be written as:

$$\mathfrak{C}_{\psi}(a,b) = \frac{1}{\sqrt{|a|}} \int_{-\infty}^{\infty} f(t) \overline{\psi\left(\frac{t-b}{a}\right)} dt, \qquad (1.4.5)$$

where "a" is called the scale factor and represents the inverse of the signal frequency, "b" is a time translation term. The mother wavelet function ϕ is continuous and differentiable with compact support. On discretizing the parameters a and bpresented in above formula one can also define the discrete wavelet transform. For that let $a = a_0^j$ where $a_0 > 1$ and $j \in \mathbb{N}$ and let $b = kb_0a_0^j$, where $k \in \mathbb{N}$ and $b_0 > 0$. In particular, $a_0 = 2$ and $b_0 = 1$ gives the following family of wavelet indexed in \mathbb{Z} commonly known as dyadic wavelet transform,

$$\psi_{j,k} = 2^{-j/2} \psi \left(2^{-j} t - k \right).$$

In this case, j = 1, 2, ..., n, where n is the base 2 logarithm of the number of the points forming the signal and $k = 1, 2..., 2^{j-1}$. $\{\psi_{j,k}(t)\}$ forms orthonormal basis for $L^2(\mathbb{R})$, i.e. Any function f belonging to $L^2(\mathbb{R})$ is written as:

$$f(t) = \sum_{j,k} a_{j,k} \psi_{j,k}(t),$$

where $a_{jk} = \langle f, \psi_{j,k} \rangle$.

Definition 1.10. [105] The dyadic discrete wavelet transform is written as:

$$C_{j,k} = 2^{-j/2} \int_{-\infty}^{\infty} f(t)\psi\left(2^{-j}t - k\right) dt, \qquad (1.4.6)$$

where j is the decomposition level (or scale) and k is the time lag. The discrete wavelet transform is faster than the continuous variant and still permits an accurate reconstruction of the input signal by inverse transformation .

Definition 1.11. [105] " The inverse wavelet transform is given by

$$f(t) = \frac{1}{C_{\psi}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \mathfrak{C}_{\psi}(a, b) \psi_{a,b}(t) \frac{dadb}{a^2}, \qquad (1.4.7)$$

where

$$C_{\psi} = \int_0^\infty \frac{|\hat{\psi}(\omega)|^2}{|\omega|} d\omega.$$

Various properties of wavelet transform

Let ψ_1 and ψ_2 be wavelets and let $f_1, f_2 \in L^2(\mathbb{R})$. Then:

a. Linearity:

$$\mathfrak{C}_{\psi}(a,b)[\alpha_1 f_1(t) + \alpha_2 f_2(t)] = \alpha_1 \mathfrak{C}_{\psi}(a,b)[f_1(t)] + \alpha_2 \mathfrak{C}_{\psi}(a,b)[f_2(t)],$$

where $\alpha_1, \alpha_2 \in \mathbb{R}$.

b. Time shift:

$$\mathfrak{C}_{\psi}(a,b)(T_{\beta}f(t)) = \mathfrak{C}_{\psi}(a,b-\beta)f(t),$$

where T_{β} is translation operator defined by $T_{\beta}f(t) = f(t-c)$.

c. Scaling/Dilation:

$$\mathfrak{C}_{\psi}(a,b)(D_d f(t)) = \frac{1}{\sqrt{d}} \mathfrak{C}_{\psi}\left(\frac{a}{d}, \frac{b}{d}\right), d > 0.$$

where D_d is a dilation operator defined by $D_c f(t) = \frac{1}{d} f(\frac{t}{d}), d > 0.$

d. Symmetry

$$\mathfrak{C}_{\psi}(a,b)f(t) = \overline{\mathfrak{C}_{\psi}(\frac{1}{a},\frac{-b}{a})}, a \neq 0.$$

e. Parity:

$$C_{P\psi}(a,b)pf(t) = \mathfrak{C}(a,-b)f(t),$$

where P is the parity operator defined as Pf(t) = f(-t).

f. Anti-linearity

$$\mathfrak{C}_{(\alpha_1\psi+\alpha_2\phi)}(a,b)f(t) = \overline{\alpha_1}\mathfrak{C}_{\psi}(a,b)f(t) + \overline{\alpha_2}\mathfrak{C}_{\phi}(a,b)f(t).$$

g.
$$\mathfrak{C}_{T_{\beta}\psi}(a,b)f(t) = \mathfrak{C}_{\psi}(a,b+\beta a)f(t).$$

h.
$$\mathfrak{C}_{D_d\psi}(a,b)f(t) = \frac{1}{\sqrt{d}}\mathfrak{C}_{\psi}(ad,b)f(t), d > 0$$

1.5 Some Examples of Wavelets

The best known and most used families of wavelets are the wavelets of Daubechies, Haar, Morlet, Meyer, Symlets and Coiflets. Here we discuss some properties of these wavelet families. The comprehensive detail about some standard wavelets can be found in [52, 74, 28].

- a. Daubechies Wavelet: I. Daubechies holds the distinction of having developed the most elegant wavelet function, which is now the basis for wavelet applications. Wavelets belonging to the Daubechies families are continuously differentiable and have compact support. These wavelets are particularly used as basis functions in signal and image processing and make it possible to represent functions at different levels of resolution. Daubechies wavelets donot have a closed analytic form, only the coefficients at the definition points are known. The compact domain of the definition of Daubechies wavelet scaling function ϕ is [0, 2N - 1] i.e ϕ is zero outside the domain while that of the corresponding wavelet function ψ is [1 - N, N], where N denotes the number of vanishing moments and 2N is the order of Daubechies wavelet or number of Daubechies filter coefficient. Two notations D_{2N} or Db_{2N} are used to refer to Daubechies wavelets. These wavelets are not symmetric except for N = 1, which is the Haar wavelet. The smoothness of these wavelets increases with the value of N. The father and mother wavelets of the first few Daubechies wavelet are shown in the Figure 1.2-1.4.
- b. Haar Wavelet: A Haar wavelet is the simplest type of wavelet. The Haar wavelet transform serves as the prototype for all other wavelet transforms. In discrete form, Haar wavelets are related to a mathematical operation called the Haar wavelet transform. It decomposes a discrete signal into sub-signals of half its length. It is conceptually simple, computationally fast, and is memory efficient. Another advantage is that it is exactly reversible without the edge effects. Haar wavelet has certain mathematical properties such as

compact support, orthogonality, and form basis for $L^2(\mathbb{R})$ – space. The basis is formed by translation and dilation of a scaling function and a mother wavelet. Haar wavelet is Daubechies wavelet of order 1, which admits zero vanishing moment and has a discontinuity. The graph of the scaling function and mother wavelet are shown in Figure 1.1. Also, rigorous mathematical detail is given in the section 2.2.

c. Morlet Wavelet: The mother wavelet for the Morlet wavelets family exists in both real and complex forms. Among them the most widely used complex Morlet wavelet is:

$$\psi_M(t) = e^{i\omega_0 t} e^{\frac{-t^2}{2\sigma_0^2}} + \epsilon(t),$$

and the Fourier transform of $\psi_M(t)$ is :

$$\hat{\psi}_M(t) = \sigma_0 e^{\frac{[(\omega - \omega_0)\sigma_0]^2}{2}} + \hat{\epsilon}(t),$$

where ω_0 represent the modulation frequency, and σ_0 denotes the width of the Gaussian. The first term alone in function $\psi_M(t)$ defined above does not satisfy the admissibility condition therefore, an additional term $\epsilon(t)$ has been added. However for large ω_0 (i.e. $\omega_0 > 5$), this additional term, which is also of Gaussian type, is numerically negligible (less than 10^{-5}) and will therefore, can be ignored in practice. Unlike the Haar wavelets family, the Morlet wavelet family is not orthogonal and does not have compact support. Also, their maximum energy lies within a narrow band around the origin. Two salient properties of Morlet wavelets are their symmetric nature and has a closed explicit mathematical expression. This wavelet is used in signal analysis, mechanical fault diagnosis, and ecological pattern detection.

d. Meyer Wavelet: This wavelet was first introduced by a French mathematician, Yves Meyer. It produced an orthogonal wavelet family with infinite support, and infinitely differentiable. The closed form expressions of Meyer scaling function $M(\omega)$ in the Fourier domain is:

$$\hat{M}(\omega) = \begin{cases} \frac{1}{\sqrt{2\pi}} & if \ |\omega| \le \frac{2\pi}{3}, \\ \frac{1}{\sqrt{2\pi}} \cos(\frac{\pi}{2}\nu(\frac{2|\omega|}{2\pi} - 1)) & if \ \frac{2\pi}{3} \le |\omega| \le \frac{4\pi}{3}, \\ 0 & otherwise. \end{cases}$$
(1.5.1)

Consequently, the closed from expression for Meyer wavelet function in Fourier domain is

$$\hat{\mathcal{W}}(\omega) = \begin{cases} 0 & \text{if } 0 < |\omega| \leq \frac{2\pi}{3}, \\ \frac{1}{\sqrt{2\pi}} \sin(\frac{\pi}{2}\nu(\frac{3|\omega|}{2\pi} - 1))e^{-i\frac{\omega}{2}} & \text{if } \frac{2\pi}{3} < |\omega| \leq \frac{4\pi}{3}, \\ \frac{1}{\sqrt{2\pi}} \cos(\frac{\pi}{2}\nu(\frac{\pi}{2}(\frac{3|\omega|}{2\pi} - 1))e^{-i\frac{\omega}{2}} & \text{if } \frac{4\pi}{3} < |\omega| \leq \frac{8\pi}{3}, \\ 0 & \text{otherwise}, \end{cases}$$
(1.5.2)

where auxiliary function ν is smooth function given by

$$\nu(x) = \begin{cases} 0 & if \ x < 0, \\ x & if \ 0 < x < 1, \\ 1 & if \ x > 1. \end{cases}$$
(1.5.3)

By choosing different auxiliary functions we can have variants of Meyer

wavelet. Furthermore, \hat{M} and \hat{W} are compactly supported continuous functions in the frequency domain.

- e. Symlets Wavelet: Symlets are a modified version of Daubechies wavelet proposed by the Daubechies and are more symmetric than Daubechies wavelet. Its symmetric property is useful in reducing edge effects. The family of symlets are compactly supported and from an orthogonal set. A symlet father wavelet and corresponding mother wavelet are shown in Figure 1.7.
- f. Coiflets Wavelet: On request of the Coifman, Daubechies designed Coifman wavelet and hence it is named as Coiflet. Coiflets are also orthogonal like Daubechies but their mother wavelets have the additional property of vanishing moments. This property makes the Coiflets very attractive for quadrature formulas. Coiflets are near symmetric wavelets and are mostly used in image processing. A scaling and mother Coiflet is shown in Figure 1.8.
- g. Wavelet generated by using orthogonal polynomials: Several orthogonal polynomials have also been used to construct wavelets families such as Chebyshev, Hermite, Lagender, Bernoulli, Laguerre, Jacobi, Gegenbauer, Lucas, Tylor, Chelyshkov, Genocchi wavelet.

1.6 Some Applications of wavelets

The wavelet method was created in the 19th century to improve signal analysis and processing. The main objective of this technique is to describe signals produced in the real world to characterize, identify, compress, filter, transmit, and predict. Long



FIGURE 1.1: (a) Haar scaling function, (b) Haar Mother wavelet.



FIGURE 1.2: (a) Daubechies(Db2) father wavelet, (b) Daubechies(Db2) Mother wavelet.



FIGURE 1.3: (a) Daubechies(Db3) father wavelet, (b) Daubechies(Db3) Mother wavelet.


FIGURE 1.4: (a) Daubechies(Db4) father wavelet, (b) Daubechies(Db4) Mother wavelet.



FIGURE 1.5: Real and Imaginary part of complex Morlet wavelet



FIGURE 1.6: (a) Meyer father wavelet, (b) Meyer Mother wavelet.



FIGURE 1.7: (a) Symlet father wavelet, (b) Symlet Mother wavelet.



FIGURE 1.8: (a) Coiflet father wavelet, (b) Coiflet Mother wavelet.

before this technique, the Fourier transformation (FT) (1822) was used for this action. FT decomposes the signal into a set of so-called "base" signals, cosine, and sine. However, it turned out that this method did not allow to have decomposition satisfactory. Wavelet has many applications in physics, engineering, mathematics, and other applied sciences. In this section, we have discussed a few applications of wavelet transform in several fields.

a. Applications in physics: Applications in physics: Wavelet has been widely used in developing new generation music synthesisers, detecting formants in speech analysis, and studying an underwater acoustic wave train. Geophysics and astrophysics: It has been applied in the detection and analysis of micro-earthquakes in oil exploration, analysis of local gravitational field in gravimetry, seismology, geomagnetism, astronomy, paleo-climatology etc. It has been used to identify coherent structures in turbulent fluids, diffusionlimited aggregates, and tree growth phenomena. It has applications in atomic physics where wavelet transforms are applied to analyse the higher order harmonics generated during laser-atom interactions. Wavelet transform has been used as an efficient tool in NMR spectroscopy for noise filtering and spectral line subtraction. A spectacular recent application of wavelet was the detection of gravitational waves by the LIGO and VIRGO experiments. The LIGO setup comprises two interferometers located in Hanford, WA, and Livingstone, LA, respectively, while the VIRGO interferometer located in Santo Stefano a Macerata, Pisa, Italy. Each captured a signal that represented the gravitational waves generated during the last femtosecond of the coalescence of two black holes. Though anticipated, but no one has ever observed the collision of two black holes. The presence of these gravitational waves is verified by applying non-stationary data analysis via the Morlet wavelet. This development in physics has confirmed various predictions of general relativity. More recently, the LIGO-VIRGO detector observed simultaneous emission of gravitational waves and electromagnetic waves from the two neutron stars of the spiral rotating binary system. This makes it possible to locate the source and estimate its distance. Here too, the detection and modelling of the phenomena are obtained by wavelet analysis. According to the experts, these results are fundamental and open a new era in astrophysics and cosmology.

- b. Applications in medicine and biology: Among the medical and biological applications, wavelet has been utilized to analyze electrocardiograms (for example, in predicting, detecting, and classifying atrial fibrillation), electroencephalograms, brain tumour detection, and detection of precursors of abnormalities (for example, epilepsy). Long-range correlations in DNA sequences are also studied using the wavelet transform. In Nuclear Magnetic Resonance Spectroscopy (NMR) and wavelet transforms are commonly employed to estimate spectral lines.
- c. Industrial applications: In industrial applications, wavelet has been used to monitor nuclear and electrical power plants. A nuclear power plant (NPP) is a complicated system consisting of large number of components with different physical behaviour. NPP emits different signals under continuously varying operational conditions. These signals carry valuable information which has to be analyzed for the safety of the plant. In order to extract information from these signals, wavelet theory is applied for signal processing. Another application of wavelet is the object shape recognition system, which plays a crucial role in robotics and automobile industry.
- d. Applications in mathematics: Wavelet transform is considered an efficient mathematical tool to detect and analyze the singularities and irregular structure, estimating the effective Holder exponent in solving differential and integral equations of integer and fractional order, approximation theory, operator theory and inverse problems. Numerical methods based on wavelet

theory has progressed in a variety of areas over the last two decades. Based on the construction algorithm, some wavelet-based techniques are; the wavelet weighted residual method, wavelet finite element method, wavelet boundary element, wavelet mesh-less method etc.

- e. Applications in crop science and vegetation: The wavelet technique has been widely utilised to examine the time series of vegetation. Several techniques have been developed to improve forecasting quality and detecting land change as well as crop phenology. Wavelet transform plays an efficient role in the processing of remotely sensed data. Remote sensing is a technique which uses sensors to collect information about objects (mostly the Earth's surface) without making physical contact with them. This data is contaminated and is very difficult to analyse directly. Therefore, processing the data, which includes compression, noise reduction, classification, and feature extraction, is needed and could be efficiently done using wavelet transform.
- f. Applications in artificial intelligence: Many artificial intelligence techniques are coupled with wavelet transform to build hybrid AI models to predict essential processes in hydrology such as; estimation and prediction of precipitation, forecasting of stream-flow caused by rainfall-runoff process, rainfall-runoff modelling, prediction of suspended sediment load due to turbulence in water bodies and many more hydro-climatologic applications. Wavelet has also been used in many artificial techniques for the texture identification of machine surfaces. Texture analysis models play a very important

role in manufacturing industries. It is used to study the texture and evaluate the component's roughness without touching it. The role of the wavelet in this modelling is to pre-process the subdivided images by converting them into grey scale using continuous 2D wavelet transform. Several statistical features (Mean, Median, Maximum, Minimum, Range, Standard Deviation, Permutation Entropy, Energy, Shannon Entropy, L1 norm, L2 norm, Maximum norm, Maximum Energy to Entropy Ratio, Log Energy Entropy, Sure Entropy, Threshold Entropy and Maximum relative Energy) are calculated from wavelet coefficients. Wavelet transform has been coupled with an artificial neural network to develop an intelligent model for electricity demand predictions. The theory of wavelet has its application in structural health monitoring techniques. In this technique, non-destructive sensor technology is used to detect defects and degradation in the structure. The signal received by the sensor is processed efficiently using a wavelet from which the health status of the structure can be viewed.

1.7 Delay Differential Equations

Generally, researchers construct mathematical models which are governed by differential equations in which the present state of the system depends only on the current value of the dependent variable and/or its derivative. Sometimes these models show severe inconsistency with reality, especially in real-time modelling, economics model, cell growth model and analysis of stock marketing. In order to improve the dynamics of such a mathematical model, researchers incorporated delay terms in the governing differential equation, which results in a delay differential equation. Delay differential equations (DDEs) form a special class of differential equations in which the rate of the solution, depends on the present as well as some previous value of the dependent variable and/or their derivative. DDEs are widely used to model physical sciences, biosciences, engineering, electrodynamics and economics processes. In the monograph [92], F.A Rihan analyses the qualitative and quantitative features of DDEs along with their applications in biosciences, [48] studies the human neural balance control model using DDEs, [110] produced some results on the convergence of Nicholson's blowflies delay model. [109] used the DDEs model to study the influence of rainfall on cocoa yield at the farm level. Karatza & Karahis [54] actively use the DDEs to develop a population pharmacokinetic model, which is an appropriate approach to describe dual peaks in irbesartan's concentration-time profiles. Another important application of DDEs is in electrohydraulic servomechanisms (EHS); this mechanism is widely used in control systems as actuators [44]. The detailed analysis and applications of DDEs can be found in [13, 35, 56, 101] and the following chapters.

Depending upon the nature of delay/lag τ , DDEs have various formats such as

• DDEs with constant delay

$$\begin{cases} y'(t) = f(t, y(t), y(t - \tau)), t_0 \le t \le t_f, \\ y(t) = \phi_0(t), t \le t_0. \end{cases}$$
(1.7.1)

• DDEs with time dependent delay $\tau(t)$

$$\begin{cases} y'(t) = f(t, y(t), y(t - \tau(t))), t_0 \le t \le t_f, \\ y(t) = \phi_0(t), t \le t_0. \end{cases}$$
(1.7.2)

• DDEs with state dependent delay $\tau(t, y)$

$$\begin{cases} y'(t) = f(t, y(t), y(t - \tau(t, y))), t_0 \le t \le t_f, \\ y(t) = \phi_0(t), t \le t_0. \end{cases}$$
(1.7.3)

• Neutral DDEs

$$\begin{cases} y'(t) = f(t, y(t), y(t - \tau_1(t, y)), y'(t - \tau_2(t, y))), t_0 \le t \le t_f, \\ y(t) = \phi_0(t), t \le t_0. \end{cases}$$
(1.7.4)

• Proportional DDEs or Pantograph equations

$$\begin{cases} y'(t) = f(t, y(t), y(qt)), t_0 \le t \le t_f, \\ y(t) = \phi_0(t), t \le t_0. \end{cases}$$
(1.7.5)

In the above equations, unlike the ordinary differential equations where we required an initial condition to solve the initial value problem here, we need an initial function $\phi_0(t)$ which is called a history function for obtaining the unique solution. The dynamical structure of DDEs is much richer than ordinary differential equations, i.e., the oscillatory and even chaotic behaviours can occur in the scalar case. Due to the involvement of delays term, the computational complexities of these equations increase; thus, it becomes too complicated to solve these equations analytically. Also, in some cases, the analytical solution of these equations does not exist; therefore, an efficient numerical technique is necessary to find the approximate solution. Basic numerical techniques for solving DDEs originates from the techniques available for ordinary differential equations, where additional operations like the inclusion of delay term are required.

1.8 Existence and Uniqueness

Here we state some results on the existence and uniqueness of the solution of DDEs. These results are obtained as particular instances of theorems proved in [45, 78],and [35]

Theorem 1.12. [13] (Local existence) Let f(t, u, v) be a continuous function on $A \subseteq [t_0, t_f) \times \mathbb{R}^d \times \mathbb{R}^d$ and has locally bounded first derivative w.r.t u and v. Moreover, let $\tau(t) \ge 0$ be a delay function which is continuous in $[t_0, t_f), \tau(t_0) = 0$ \mathscr{C} , for some $\xi > 0$, $t - \tau(t) > t_0$ in the interval $(t_0, t_0 + \xi]$. Then equation

$$\begin{cases} y'(t) = f(t, y(t), y(t - \tau(t))), t_0 \le t \le t_f, \\ y(t_0) = y_0, \end{cases}$$
(1.8.1)

has unique solution in $[t_0, t_0 + \delta)$ for some $\delta > 0$ which depends continuously on the initial data.

Theorem 1.13. [13] (Global existence) Let f(t, u, v) be a continuous function on $A \subseteq [t_0, t_f) \times \mathbb{R}^d \times \mathbb{R}^d$ and has locally bounded first derivative w.r.t u and v. Moreover, let $\tau(t) \ge 0$ be a delay function which is continuous in $[t_0, t_f), \tau(t_0) = 0$ \mathscr{C} , for some $\xi > 0$, $t - \tau(t) > t_0$ in the interval $(t_0, t_0 + \xi]$. If the unique maximal solution of equation

$$\begin{cases} y'(t) = f(t, y(t), y(t - \tau(t))), t_0 \le t \le t_f, \\ y(t_0) = y_0, \end{cases}$$
(1.8.2)

defined in the interval $[t_0, b)$, with $t_0 < b \le t_f$ is bounded, then it exists on the entire interval $[t_0, t_f)$. In [35], Driver has proved the Theorem 1.12, and some more general cases of state-dependent delay. For further developments and more general cases, we refer interested readers to [27, 45], and the references therein.

1.9 Methods for Solving DDEs

Here we have highlighted some analytical methods for solving delay differential equations.

1.9.1 Method of Steps for Solving DDEs

Delay differential equation may be solved as ordinary differential equations over successive intervals $[t_l, t_{l+1}]$ by the method of steps. The scalar DDE

$$y'(t) = f(t, y(t), y(t - \tau)), t > 0,$$
(1.9.1)

with initial function $\phi_0(t)$ defined on $t \in [-\tau, 0)$ and initial condition $y(0) = y_0$ is solved as a chain of differential equation.

$$y_{1}'(t) = f(t, y_{1}(t), \phi_{0}(t - \tau)), \quad 0 \leq t \leq \tau,$$

$$y_{2}'(t) = f(t, y_{2}(t), y_{1}(t - \tau)), \quad \tau \leq t \leq 2\tau,$$

$$\vdots$$

$$y_{n}'(t) = f(t, y_{n}(t), y_{n-1}(t - \tau)), \quad (n - 1)\tau \leq t \leq n\tau,$$

(1.9.2)

where $n \in \mathbb{Z}^+$.

Example 1.1. consider a simple DDEs with constant delay

$$y'(t) = y(t-1) \tag{1.9.3}$$

with initial function $y(t) = 1, \forall t \in [-1, 0]$ and initial condition y(0) = 1.

Using basic idea of method of steps, first we reduce the given DDEs on the interval [0, 1] to a non autonomous ODE as

$$y_1'(t) = \phi_0(t-1)$$

where $\phi_0(t) = 1$, $t \in [-1, 0]$. Now by using integral form of the solutions, we have $\forall t \in [0, 1]$

$$y'_{1}(t) = y(0) + \int_{0}^{t} \phi_{0}(s-1)ds$$

= $y(0) + \int_{0}^{t} 1ds$
= $y(0) + t$

= 1 + t

Now y(t) (denoted by $y_1(t)$) is known in [0,1], proceeding as before we can reduce the DDE (1.9.3) on the interval [1,2] to a non-autonomous ODE as

$$y_2'(t) = y_1(t-1)$$

where $y_1(t) = 1 + t$, $t \in [0, 1]$. Again by using integral form of the solution for $t \in [1, 2]$, we have

$$y_{2}'(t) = y_{1}(1) + \int_{1}^{t} y_{1}(s-1)ds$$
$$= y_{1}(1) + \int_{0}^{t} (1+t)ds$$
$$= y_{1}(1) + t + \frac{t^{2}}{2}$$
$$= 2 + t + \frac{t^{2}}{2}$$

Therefore, for $t \in [1, 2]$

$$y_2(t) = 2 + t + \frac{t^2}{2}.$$

One can continue the procedure until the desired solution is obtained. But most of the time the resulting integrals quickly become very cumbersome and it is very difficult to draw any conclusion about the solution from this exact procedure.

1.9.2 Laplace Transformation Method

Consider the equation

$$y'(t) = cy(t-1) + f(t) \quad t > 1, \tag{1.9.4}$$

with initial function y(t) = g(t), $0 \le t \le 1$. Taking the Laplace transforms on both side leads to

$$\int_{1}^{\infty} y'(t)e^{-st}dt = \int_{1}^{\infty} y'(t-1)e^{-st}dt + \int_{1}^{\infty} f(t)e^{-st}dt$$
$$y(t)e^{-st}dt - \int_{1}^{\infty} y'(t)(-s)e^{-st}dt = c\int_{0}^{\infty} y'(u)e^{-s(u+1)}du.$$

Assume that $y(t)e^{-st} \to 0$ as $t \to \infty$, this leads to

$$-e^{-s}y(1) + s\int_{1}^{\infty}y(t)e^{-st}dt = ce^{-s}\int_{0}^{1}y(u)e^{-su}du + ce^{-s}\int_{1}^{\infty}y(u)e^{-su}du$$

Also assume that $s - ce^{-s} \neq 0$, then we have

$$\int_{1}^{\infty} y(t)e^{-st}dt = \frac{ce^{-s}\int_{0}^{1} y(u)e^{-su}du + e^{-s}y(1) + \int_{1}^{\infty} f(t)e^{-st}dt}{s - ce^{-s}}.$$
 (1.9.5)

Assume that the inversion of the Laplace transform can be applied to have

$$y(t) = \int_{(b)} \left(\frac{ce^{-s} \int_0^1 y(u)e^{-su} du + e^{-s} y(1) + \int_1^\infty f(t)e^{-st} dt}{s - ce^{-s}} \right) e^{st} ds, \qquad (1.9.6)$$

where the integral is carried out over a vertical line Re(s) = b, with b sufficiently large. It turns out to be the case that all zeros of $s - ce^{-s}$ lie in a left half-plane, and that the relation in (1.9.6) is valid if all these zeros satisfy Re(s) < b. In [27], other methods such as solution by definite integrals, series expansions, distribution of characteristics Roots along with the asymptotic behaviour of the solution and problem stability are discussed.

The purpose of this thesis is to present a reliable numerical approximation of a special class of delay differential equations known as Pantograph equation. The simplest form of proportional delay differential equation is

$$y'(t) = ry(t) + sy(qt), \ 0 < q < 1, \ t \ge 0,$$
(1.9.7)

where $r, s, \in \mathbb{C}$.

In some literature, these equations are mentioned as proportional delay differential equations. These equations explain many physical and mathematical phenomena, such as the dynamics of an overhead current collection system for an electric locomotive, light absorption in the galaxy, queuing theory, partition problems in number theory, cell growth model, and probability theory. Numerical solutions of IVPs, BVPs, a systems of differential equations, and fractional-order differential equations, along with their applications have been discussed in the following chapters.

1.10 Objective

This research formulates new numerical techniques for solving proportional delay differential equations. The Haar wavelet series method and the modified Haar wavelet series method are constructed to solve linear and nonlinear initial value problems (IVPs), boundary value problems, systems of differential equations, and fractional order IVPs. These schemes are capable of producing results with high accuracy and with less computational cost. All computer simulations are carried out in MATLAB on a 16GB RAM 64 bit machine carrying an Intel i5 processor. These methods involve fewer cumbersome manual calculations as compared to other numerical schemes.

1.11 Thesis Outline

The thesis discussed Haar wavelet techniques to solve different types of proportional delay differential equations. In Chapter 2, we introduced some preliminaries and definitions of Haar wavelets and their integrals which will be used throughout the thesis and developed a numerical technique to solve the proportional delay Riccati differential equation using Haar wavelets. Chapter 3 of the thesis deals with the Haar wavelet-based numerical technique for solving the proportional delay variants of Dirichlet boundary value problems. We have applied the Haar wavelet series method for solving simultaneous proportional delay differential equations in Chapter 4. Chapter 5 presents the modified Haar wavelet series method to solve higher-order multi-pantograph equations arising in electrodynamics. In Chapter 6, we discussed some concepts of fractional calculus and developed a reliable collocation numerical technique to solve proportional delay Riccati differential equations of Fractional Order.

Chapter 2

Approximate solution for proportional delay Riccati differential equations by Haar wavelet method

2.1 Riccati Differential Equations

The Riccati differential equations (RDEs) come under the class of nonlinear differential equations. These equations are widely studied for numerous problems of contemporary analysis and their applications and are not easy to solve explicitly. This makes it interesting to investigate the solutions of these equations. Here, we have considered the following Riccati differential equations:

$$y'(t) = q_1(t) + y(t) \left(q_2(t) + q_3(t)y(t) \right), \quad t_0 \le t \le t_f, \quad y(t_0) = y_0, \tag{2.1.1}$$

where $q_1(t), q_2(t)$ and $q_3(t) \neq 0$ are continuous, t_0, t_f and $y(t_0)$ are arbitrary constant and y(t) is unknown function.

Its proportional-delay variant can be written as

$$y'(t) = \psi(t) + by(t) + cy(\alpha t)(d - y(\alpha t)), \quad t_0 \le t \le t_f, \quad y(t_0) = y_0, \qquad (2.1.2)$$

where $c \neq 0$, $b, d, y_0 \in C$, and $\alpha > 0, \psi(t)$ is a continuous and $\alpha \neq 1$. When $0 < \alpha < 1$ equation (2.1.2) yields a retarded equation, whereas $\alpha > 1$ produces advanced equation. It follows from Picard-Lindelöf theorem that the solutions of equation (2.1.1) and equation (2.1.2) exists and it is unique. For further studies related to Riccati differential equations see [14, 75, 90, 91].

RDEs widely appear in random processes, kalman filtering systems, and network

synthesis. It has enormous applications in the fields of super symmetric quantum mechanics and quantum chemistry and plays a key role in the theory of optimal control, stochastic control theory, financial mathematics, diffusion problems, econometric models, and dynamic games.

Modeling by delay is necessary in many applied physical problems, and to accomplish the purpose, delay differential equations have been used significantly. Delay differential equations arise inevitably in decision making, mathematical modelling of chemicals, biological and physiological processes, economic growth, neural networks, and delayed dynamics. For further applications of RDEs, one can see [23, 34, 59, 76] and the references mentioned in the following chapters.

In recent years, the problem of finding the approximate solution of these differential equations has grabbed attention and has been examined by many mathematicians. The methods already used for the numerical solution of Riccati differential equations are the variational iteration method (VIM)[38], the modified variational iteration method (MVIM)[1], homotopy perturbation method (HPM)[16], differential transform method (DTM)[75], and the Bezier curves method[40]. The Reproducing kernel Hilbert space method (RKHSM)[49] and Bezier control point method [39] are used to find the approximate solution of delay RDEs. In [23], semi-analytical solutions of some nonlinear proportional delay differential equations are discussed. In [5], author studied the numerical treatment of the stochastic delay differential equation to formulate a one-step scheme to approximate the solution of the problem. [80] studied the numerical inclusion of exact periodic solutions for the delay Duffing equation and proposed the existence of periodic solutions for the forced delay Duffing equations based on the verified numerical computations.

The aim of this chapter is to present a numerical method which has a low computational cost and will allow us to solve RDEs and proportional delay RDEs. The method is based on the Haar wavelet basis and is named as Haar wavelet method (HWM). The idea is to convert a differential equation into a system of algebraic equations from which an approximate solution can be obtained. The method is tested on linear as well as nonlinear problems, and very promising results are obtained.

2.2 Haar Wavelet Basis and its Integration

To construct the Haar wavelet system $\{\mathfrak{h}_{i}(t)\}_{i=1}^{\infty}$ on $[\Gamma_{1}, \Gamma_{2}]$ two basic functions are required, namely

(a). The Haar scaling function (father wavelet):

$$\mathfrak{h}_1(t) = \mathbb{I}_{[\Gamma_1, \Gamma_2)}(t). \tag{2.2.1}$$

(b). The mother wavelet:

$$\mathfrak{h}_{2}(t) = \mathbb{I}_{[\Gamma_{1},(\Gamma_{1}+\Gamma_{2})/2)}(t) - \mathbb{I}_{[(\Gamma_{1}+\Gamma_{2})/2,\Gamma_{2})}(t), \qquad (2.2.2)$$

where $\mathbb{I}_{[a,b]}(t)$ is characteristic/indicator function. Now for generating the Haar wavelet series, let j be dilation and k be translation parameter.

Then i-th Haar Wavelet is defined as:

$$\mathfrak{h}_{i}(t) = \begin{cases} 1 & \text{for } t \in [\vartheta_{1}(i), \vartheta_{2}(i)), \\ -1 & \text{for } t \in [\vartheta_{2}(i), \vartheta_{3}(i)), \\ 0 & otherwise, \end{cases}$$
(2.2.3)

where $\vartheta_1(i) = \Gamma_1 + (\Gamma_2 - \Gamma_1)k/2^i, \vartheta_2(i) = \Gamma_1 + (\Gamma_2 - \Gamma_1)(k + 0.5)/2^i, \vartheta_3(i) = \Gamma_1 + (\Gamma_2 - \Gamma_1)(k + 1)/2^i$. The index $i = 2^i + k + 1, j = 0, 1, ..., J$ where J is maximum level of wavelet and $k = 0, 1, ..., 2^i - 1$.

(c). Define scaling function space and wavelet space as follows

$$V_{j} = span\{2^{j/2}\mathfrak{h}_{1}(2^{j}t - k), t \in [\Gamma_{1}, \Gamma_{2}]\}_{k=0}^{2^{j}-1},$$

$$W_{j} = span\{2^{j/2}\mathfrak{h}_{2}(2^{j}t - k), t \in [\Gamma_{1}, \Gamma_{2}]\}_{k=0}^{2^{j}-1}.$$
(2.2.4)

Suppose $0 \leq J_0 < J$, then following relation holds:

$$V_J = V_{J_0} \otimes W_{J_0} \otimes W_{J_0+1} \cdots \otimes W_{J-1}.$$

$$(2.2.5)$$

The spaces V_j are such that $V_0 \subset V_1 \subset V_2 \cdots \subset L^2([\Gamma_1, \Gamma_2])$ and $\overline{\bigcup_{j=0}^{\infty} V_j} = L^2([\Gamma_1, \Gamma_2])$. Hence $L^2([\Gamma_1, \Gamma_2]) = V_0 \bigotimes (\bigotimes_{j_0}^{\infty} W_j)$ holds. It allows us to approximate any $f \in L^2([\Gamma_1, \Gamma_2])$ with following truncated Haar series:

$$f_{approx}(t) \approx \sum_{i=1}^{2^{J+1}} a_i \mathfrak{h}_i(t).$$
(2.2.6)

In the interval $[0, 1], \vartheta_1(i), \vartheta_2(i), \vartheta_3(i)$ becomes:

$$\vartheta_1(i) = \frac{k}{2^j}, \ \vartheta_2(i) = \frac{(k+0.5)}{2^j}, \ \vartheta_3(i) = \frac{(k+1)}{2^j}.$$

To apply the Haar wavelet following integrals are required:

$$I_1\mathfrak{h}_i(t) = \int_0^t \mathfrak{h}_i(u) du, \ I_2\mathfrak{h}_i(t) = \int_0^t I_1\mathfrak{h}_i(u) du, \ \dots \ I_n\mathfrak{h}_i(t) = \int_0^t I_{(n-1)}\mathfrak{h}_i(u) du.$$

These integral can be calculated by analytic integration of (2.2.3),

$$I_{1}\mathfrak{h}_{i}(t) = \begin{cases} t - \vartheta_{1}(i) & \text{for } t \in [\vartheta_{1}(i), \vartheta_{2}(i)), \\\\ \vartheta_{3}(i) - t & \text{for } t \in [\vartheta_{2}(i), \vartheta_{3}(i)), \\\\ 0 & otherwise, \end{cases}$$
(2.2.7)

and

$$I_{2}\mathfrak{h}_{i}(t) = \begin{cases} \frac{1}{2}(t - \vartheta_{1}(i))^{2} & \text{for } t \in [\vartheta_{1}(i), \vartheta_{2}(i)), \\ \\ \frac{1}{2^{2j+2}} - \frac{1}{2}(\vartheta_{3}(i) - t)^{2} & \text{for } t \in [\vartheta_{2}(i), \vartheta_{3}(i)), \\ \\ \frac{1}{2^{2j+2}} & \text{for } t \in [\vartheta_{3}(i), 1), \\ \\ 0 & \text{otherwise.} \end{cases}$$
(2.2.8)

In general, the value of nth integral of (2.2.3) on interval [0, 1] is

$$I_{n}\mathfrak{h}_{i}(t) = \frac{1}{n!} \begin{cases} 0 & \text{for } t \in [0,\vartheta_{1}(i)), \\ (t - \vartheta_{1}(i))^{n} & \text{for } t \in [\vartheta_{1}(i),\vartheta_{2}(i)), \\ (t - \vartheta_{1}(i))^{n} - 2(t - \vartheta_{2}(i))^{n} & \text{for } t \in [\vartheta_{2}(i),\vartheta_{3}(i)), \\ (t - \vartheta_{1}(i))^{n} - 2(t - \vartheta_{2}(i))^{n} + (t - \vartheta_{3}(i))^{n} & \text{for } t \in [\vartheta_{3}(i), 1). \end{cases}$$

$$(2.2.9)$$

In case i = 1, we have the n^{th} integral of scaling function evaluated from (2.2.1)

$$I_n\mathfrak{h}_i(t) = \frac{t^n}{n!}.$$
(2.2.10)

2.3 Description of Method

In order to solve equation (2.1.1) and equation (2.1.2), by Haar wavelet expansion, let

$$y'(t) = \sum_{i=1}^{2M} a_i \mathfrak{h}_i(t).$$
 (2.3.1)

On integrating equation (2.3.1) from 0 to t with respect to t we compute

$$y(t) = \sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(t) + y(0), \qquad (2.3.2)$$

also

$$y'(\alpha t) = \sum_{i=1}^{2M} a_i \mathfrak{h}_i(\alpha t), \qquad (2.3.3)$$

and

$$y(\alpha t) = \sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(\alpha t) + y(0). \qquad (2.3.4)$$

Using equations (2.3.1), (2.3.2) in equation (2.1.1) and equations (2.3.3), (2.3.4) in equation (2.1.2) with collocation points $t_l = \frac{(l-0.5)}{2M}$, $l = 1, 2, \ldots, 2M$, we have the following system of equations

$$\sum_{i=1}^{2M} a_i \mathfrak{h}_i(t_l) = q_1(t_l) + \left[\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(t_l) + y(0)\right] q_2(t_l)$$

$$+ q_3(t_l) \left[\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(t_l) + y(0)\right] \left[\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(t_l) + y(0)\right],$$
(2.3.5)

and

$$\sum_{i=1}^{2M} a_i \mathfrak{h}_i(t_l) = \psi(t_l) + b \left[\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(t_l) + y(0) \right]$$

$$+ c \left[\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(\alpha t_l) + y(0) \right] \left[d - \left(\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(\alpha t_l) + y(0) \right) \right].$$
(2.3.6)

Above systems of nonlinear equations can be solved using classical "Newton's method", by using these values of Haar coefficient a_i 's obtained from equation (2.3.5) in equation (2.3.2), we get the approximate solution of equation (2.1.1), similarly using coefficients obtained form equation (2.3.6) in equation (2.3.2), we get the approximate solution of equation (2.3.2), we

2.3.1 Convergence analysis of the Haar wavelet

Lemma 2.1. Assume that $\omega(t) \in L^2(\mathbb{R})$ be a Lipschitz continuous (with $K = \sup |\omega'(t)|$) on [0, 1), then the error norm at Jth level satisfies the following inequality

$$\|e_j(t)\| \le \sqrt{\frac{K}{7}} \frac{C}{2^{(3)2^{J-1}}},\tag{2.3.7}$$

where K, C are some real constants.

Proof. For proof see [9].

2.4 Numerical Problems

Here we will present numerical problems and the discussion on the Haar wavelet method for solving different types of Riccati and proportional delay Riccati differential equations. We have shown that the performance of our technique is sharp and the absolute error is reduced significantly in some problems. We also calculated the experimental rate of convergence $R_c(J)$, as described in [68]. In the first part of this section, we consider quadratic Riccati differential equations. Then, we solve some delay-Riccati differential equations.

$$R_c(J) = \frac{\log(E_c(J-1)/E_c(J))}{\log(2)},$$
(2.4.1)

where $E_c(J)$ is error at level J.

Problem 1. Consider the following, taken from [40]

$$y'(t) = 16t^2 - 5 + 8ty(t) + y^2(t), \quad 0 \le t \le 1, \quad y(0) = 1.$$
 (2.4.2)

The exact solution of equation (2.4.2) is y(t) = 1 - 4t. Here we have solved this problem using Haar wavelet expansions and integral of the Haar wavelet. The values

Value of J	\mathbf{MAEs}	
2	1.4654E - 14	
3	1.1546E - 14	
4	3.1086E - 15	
5	2.6645E - 15	
6	3.5527E - 15	
7	8.8818E - 16	

TABLE 2.1: MAEs at different level of J (Problem 1).

TABLE 2.2: MAEs for Problem 1 with J=3.

HWM	Bezier Curves Method[40]
5.3735E - 14	5.3714E - 04

of Haar coefficients at level J = 2 are $a_1 = -4.000000, a_2 = 0.000000, a_3 = 0.000000, a_4 = 0.000000, a_5 = 0.000000, a_6 = 0.000000, a_7 = 0.000000, a_8 = 0.000000$. The coefficient a_1 is significantly close to -4 and rest of the coefficients are zero correct up to six decimal places. Therefore, from equation (2.3.2) solution is $a_1I_1h_1(t) + y(0)$ i.e.-4t + 1 which is same as exact solution. The error estimates at different levels of resolution are given in Table 2.1. We observed that good approximation can be achieve with the increase in level of resolution. A comparison between the Haar wavelet method and the Bezier Curves method [40] is reported in Table 2.2, which clearly shows that the absolute error is reduced significantly in our case as compared to the Bezier Curves method. Point wise error at J = 3 is shown in Table 2.3. Comparison between exact solution and numerical solution at J = 3 is shown in Fig.2.1 and we observed that both the curves visually coincide.

Problem 2. Solve the equation

$$y'(t) = e^{t} - e^{3t} + 2e^{2t}y(t) - e^{t}y^{2}(t), \quad 0 \le t \le 1, \quad y(0) = 1.$$
(2.4.3)



FIGURE 2.1: Comparison of exact and numerical solution at J = 3.

$t_l = (l - 0.5)/2^{J+1}$			
t(=1/32)	Haar solution	\mathbf{Exact}	Error
1	0.8750	0.8750	0.0222E - 014
3	0.6250	0.6250	0.0666E - 014
5	0.3750	0.3750	0.0666E - 014
7	0.1250	0.1250	0.0333E - 014
9	-0.1250	-0.1250	0.0222E - 014
11	-0.3750	-0.3750	0.0888E - 014
13	-0.6250	-0.6250	0.1554E - 014
15	-0.8750	-0.8750	0.1554E - 014
17	-1.1250	-1.1250	0.1332E - 014
19	-1.3750	-1.3750	0.1332E - 014
21	-1.6250	-1.6250	0.1332E - 014
23	-1.8750	-1.8750	0.0444E - 014
25	-2.1250	-2.1250	0.0888E - 014
27	-2.3750	-2.3750	0.2665E - 014
29	-2.6250	-2.6250	0.4885E - 014
31	-2.8750	-2.8750	0.8882E - 014

TABLE 2.3: Comparison between exact solution and Haar solution (Problem 1).

The exact solution of equation (2.4.3) is $y(t) = e^t[40]$. The maximum absolute error at different levels of resolution is shown in Table 2.4. From Table 2.4, we have observed that error is decreasing with the increase in resolution. The result is compared with the Bezier Curves method [40] in Table 2.6. A comparison between the exact solution and the numerical solution for J = 4 is depicted in Fig.2.2.



FIGURE 2.2: Comparison between exact and numerical solution (Problem 2).

TABLE 2.4:	MAEs at different levels
of	J (Problem 2).

TABLE 2.5: MAEs at different levels of J (Problem 4).

Value of J	MAEs	Value of J	MAEs
4	2.5841E - 04	5	3.1062E - 03
5	6.5049E - 05	6	7.8308E - 04
6	1.6315E - 05	7	1.9648E - 04
7	4.0855E - 06	8	4.9204E - 05
8	1.0238E - 06	9	1.2307E - 05
9	2.5889E - 07	10	3.0790E - 06

TABLE 2.6: Comparison for Problem 2.

HWM	Bezier Curves Method[40]
6.5049E - 05	8.0175E - 04

Computer simulations indicate that by increasing resolution, a good approximate solution can be achieved.

Problem 3. In this example, we choose to solve the following Riccati differential equation

$$y'(t) = y(t) - 50ty^2 + f(t), \quad 0 \le t \le 1, \quad y(0) = 1.$$
 (2.4.4)

The function f(t) is chosen such that equation (2.4.4) has exact solution y(t) =



FIGURE 2.3: Comparison between exact and numerical solution (Problem 3).

 $\frac{1}{1+25t^2}$. We have calculated the maximum absolute error and experimental rate of convergence at J = 4, 5, 6, 7, 8, 9. The maximum absolute error is decreasing from order 10^{-3} for J = 4 to order 10^{-6} for J = 9. Further, we observed that the maximum absolute error decreases by increasing the levels of resolution and the numerical rate of convergence approaches 2, thus confirming the theoretical results (studied by the authors in [68]). Finally, the comparison between approximate and exact solution is depicted in Fig.2.3

Problem 4. In this illustration we consider the following PDRDE

$$y'(t) = 1 - 2y^2(\frac{t}{2}), \quad 0 \le t \le 2\pi, \quad y(0) = 0.$$
 (2.4.5)

The exact solution of equation (2.4.5) is y(t) = sin(t)[23]. In this problem $\alpha = 1/2$, which is the case of a retarded equation. We have solved this problem using Haar wavelet expansion and integral of the Haar wavelet. The maximum absolute error for different resolutions is shown in Table 2.5. A comparison between analytic and approximate solutions at J = 5 is shown in Fig.2.4. The maximum absolute error is 3.0790E - 06 at J = 10. We observed that by increasing the value of J a good



FIGURE 2.4: Comparison between exact and numerical solution (Problem 4).

approximation can be achieved.

Problem 5. Solve the delay equation

$$y'(t) = \frac{1}{4}y(t) + y(\frac{t}{2})(1 - y(\frac{t}{2})), \quad 0 \le t \le 1, \quad y(0) = 1.$$
 (2.4.6)

The equation equation (2.4.6) is an example of proportional-delay variant of Riccati differential equation and it possesses the periodic solution

$$y(t) = \frac{1}{2} + \frac{1}{2}\cos(\frac{\sqrt{2}t}{4}) + \frac{\sqrt{2}}{2}\sin(\frac{\sqrt{2}t}{4}).$$

First, we will transform equation (2.4.6) into the following system of algebraic equations,

$$\sum_{i=1}^{2M} a_i \mathfrak{h}_i(t_l) - \frac{1}{4} \left[\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(t_l) + 1 \right] - \left[\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(\frac{t_l}{2}) + 1 \right] \left[1 - \left(\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(\frac{t_l}{2}) + 1 \right) \right] = 0$$

$$(2.4.7)$$

By solving the system of algebraic equation (2.4.7), we obtained the values of the Haar coefficients a_i . Using the values of a_i along with given initial condition y(0) = 1 in equation (2.3.2) approximate numerical solution for equation (2.4.6) is obtained. Maximum absolute error for different values of J is presented in Table

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MAEs (Problem 5).	TABLE 2.8 :	MAEs (Problem 6).
MAEs	Value of J	MAEs
3.0352E - 05	4	5.8484E - 04
7.6091E - 06	5	1.5242E - 04
1.9048E - 06	6	3.8914E - 05
4.7652E - 07	7	9.8315E - 06
1.1917E - 07	8	2.4709E - 06
7.4500E - 09	9	6.1935E - 07
	MAEs (Problem 5). MAEs $3.0352E - 05$ 7.6091E - 06 1.9048E - 06 4.7652E - 07 1.1917E - 07 7.4500E - 09	MAEs (Problem 5).TABLE 2.8:MAEsValue of J $3.0352E - 05$ 4 $7.6091E - 06$ 5 $1.9048E - 06$ 6 $4.7652E - 07$ 7 $1.1917E - 07$ 8 $7.4500E - 09$ 9



2.7 and the maximum absolute error is 7.4500E - 09 for J = 8. In Table 2.9 point wise error between exact solution and Haar solution at J = 3 is reported. Finally, comparison between exact solution and numerical solution at J = 3 is depicted graphically in Fig.2.7. We observed that a better approximation can be achieved with an increase in the value of J.

Problem 6. Solve the delay equation

$$y'(t) = -\frac{1}{8}y(t) + y(\frac{t}{2})(1 - y(\frac{t}{2})), \quad 0 \le t \le 4\pi, \quad y(0) = \frac{1}{4}.$$
 (2.4.8)



FIGURE 2.7: Comparison between exact and numerical solution (Problem 5).

$t_l = (l - 0.5)/2^{J+1}$			
t(=1/32)	Haar solution	Exact solution	Error
1	1.007751	1.007781	0.303515E - 004
3	1.023128	1.023158	0.300143E - 004
5	1.038250	1.038279	0.296578E - 004
7	1.053109	1.053138	0.292823E - 004
9	1.067697	1.067726	0.288881E - 004
11	1.082009	1.082038	0.284758E - 004
13	1.096036	1.096065	0.280460E - 004
15	1.109773	1.109801	0.275987E - 004
17	1.123212	1.123239	0.271345E - 004
19	1.136346	1.136373	0.266540E - 004
21	1.149170	1.149196	0.261577E - 004
23	1.161677	1.161702	0.256458E - 004
25	1.173860	1.173885	0.251189E - 004
27	1.185715	1.185740	0.245777E - 004
29	1.197235	1.197259	0.240227E - 004
31	1.208414	1.208438	0.234542E - 004

TABLE 2.9: Comparison between exact solution and Haar solution (Problem 5).

Note that equation (2.4.8) is first order Riccati differential equation with proportional delay. The exact solution for this problem is

$$y(t) = \frac{1}{2} - \frac{1}{4}\cos(\frac{\sqrt{5}t}{8}) + \frac{\sqrt{5}}{4}\sin(\frac{\sqrt{5}t}{8}).$$

On applying the numerical technique mentioned in section 3, the maximum absolute



error between the exact solution and the numerical solution at different values of J is recorded in Table 2.8. From Table 2.8 we have observed that the maximum absolute error decreases with the increase in resolution. The maximum absolute error is decreased from the order of 10^{-4} at J = 4 to 10^{-7} at J = 9 and the rate of convergence approaches 2, which follows the theoretical results mentioned in [68]. A comparison between the exact solution and the numerical solution at J = 4 is shown in Fig.2.8.

2.5 Summary

In this chapter¹, the Haar wavelet method has been successfully applied to find the numerical solution of Riccati differential equations and proportional-delay variants of Riccati differential equations, and several problems from the literature have been solved. Numerical results are compared with the Bezier Curves method and exact solutions. The numerical simulations show that our computation is much better in terms of accuracy and convergence.

¹The content of this chapter is published in **Poincare Journal of Analysis and Applica**tions .(Scopus Indexed).

Chapter 3

Haar based numerical technique for solving proportional delay Dirichlet boundary value problems

3.1 Introduction

Boundary value problems(BVPs) are now prevalent in all applied sciences. The growing popularity of this branch of differential equations has prompted numerous researchers to investigate various physical models utilizing mathematical tools and computer simulation software. Many problems in science and technology such as study of the mechanical behavior of the nano material in nanomechanics[42], turbulence modeling[58], modeling of chemical reactors[47], study of molecular structure in chemical engineering[12], heat transfer model and study of deflection in cables can be formulated mathematically in BVPs for second order differential equations.

Another application of BVPs appears in mathematical modelling design to reflect the real mechanical properties of smart material[69]. The smart material, also called as intellectual material of the 21st century, is a material with shape memory effects. Such materials are widely used in medicine, engineering, aircraft building, construction, etc.

The main focus of this study is to develop a Haar wavelet-based numerical technique for solving the following proportional delay variant of the two-point boundary value problem

$$y''(t) = \Omega(t, y(t), y(qt), y'(t), y'(qt)), \ t \in [0, 1],$$

with $y(0) = \zeta_1, \ y(1) = \zeta_2,$ (3.1.1)

where Ω is analytic function, $q \in (0, 1)$ and, ζ_1, ζ_2 are arbitrary constants.

The proportional delay variant of the delay differential equation is generally referred to as the Pantograph equation. The name Pantograph comes from Ockendon and Tayler's first work for collecting current by the Pantograph head of an electric locomotive. These equations are used in many different fields, such as number theory, probability theory on algebraic structure, economics, cell growth modelling, astrophysics, nonlinear dynamical systems, adaptive control, quantum mechanics, electrodynamics, engineering, and numerous others. [13, 35, 36, 111].

In many instances, it is difficult to obtain the exact solution of the delay differential models. As a result, the solutions to such equations have developed a lot of interest among researchers, and they have used a variety of numerical approaches to approximate the solutions to these equations. Agarwal and Chow extended the finite-difference method to offer the approximate solution to two-point boundaryvalue problems with deviating argument [3]. Li and Liu presented a novel numerical approach for solving multi-pantograph delay equations based on the Runge-Kutta scheme [61]. Shakeri et al. used the homotopy perturbation method (HPM) to solve certain delay differential equations [100]. To solve generalised pantograph equations, Saadatmandi et al. used a Lagrange multiplier-based variational iteration scheme (VIM) [93]. Both VIM and HPM yield correct results, but the computational cost is significant due to the use of symbolic integrations. Shakeri and Dehghan have investigated the numerical solution of delay differential equations using the adomain decomposition method [99]. Sedaghat et al. [96] adopted a Chebyshev polynomial based numerical method and approximate the solution of DDEs. The aim of the method [96] is to utilise the operational matrix and its derivative to simplify the problem into a series of algebraic equations from which the solution can be derived. Authors in [15, 33, 112, 115] have developed some numerical techniques based on collocation method in the recent past. These techniques are quite effective for investigating proportional delay differential equations.

The author is inspired by the aforementioned findings and adopted the Haar wavelet series approach for examining the approximate numerical solutions of Dirichlet boundary value problems of proportional delay nature. The method is suitable for solving BVPs since it takes the boundary condition into consideration automatically. Moreover, for a limited number of collocation points, high precision is achievable.

3.2 Construction of Method

We consider the boundary value problem of pantograph type as follows:

$$y''(t) = \Omega(t, y(t), y(qt), y'(t), y'(qt)), \ t \in [0, 1],$$

(3.2.1)
with Dirichlet condition $y(0) = \zeta_1, \ y(1) = \zeta_2,$

where Ω is analytic function, $q \in (0, 1)$ and, ζ_1, ζ_2 are arbitrary constant. In order to apply the Haar wavelet series techniques, firstly we expand the y''(t) in terms of truncated Haar wavelet series as:

$$y''(t) = \sum_{i=0}^{2^{J+1}} a_i \mathfrak{h}_i(t).$$
(3.2.2)

Integrate equation (3.2.2) 0 to t, we get,

$$y'(t) = \sum_{i=0}^{2^{J+1}} a_i I_1 \mathfrak{h}_i(t) + y'(0).$$
(3.2.3)

Further integration yields,

$$y(t) = \sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(t) + y'(0)t + y(0).$$
 (3.2.4)

Quantity y'(0) in equation (3.2.3) and (3.2.4) are yet to determined for that integrate equation (3.2.3) from t to 1, we have

$$-y(t) = \sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - \sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(t) + y'(0)(1-t) - y(1).$$
(3.2.5)

Now, from equation (3.2.4) and (3.2.5) we get,

$$y'(0) = -\sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - y(0) + y(1).$$

Utilizing y'(0), $y(0) = \zeta_1$, $y(1) = \zeta_2$ in equation (3.2.4), we obtain the following

$$y(t) = \sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(t) + \left(-\sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - \zeta_1 + \zeta_2\right) t + \zeta_1.$$
(3.2.6)

Also,

$$y(qt) = \sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(qt) + \left(-\sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - \zeta_1 + \zeta_2\right)(qt) + \zeta_1.$$
(3.2.7)

Similarly, putting $y'(0) = -\sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - y(0) + y(1), \ y(0) = \zeta_1, \ y(1) = \zeta_2$ in equation (3.2.3) we get,

$$y'(t) = \sum_{i=0}^{2^{J+1}} a_i I_1 \mathfrak{h}_i(t) - \sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - \zeta_1 + \zeta_2, \qquad (3.2.8)$$

and

$$y'(qt) = \sum_{i=0}^{2^{J+1}} a_i I_1 \mathfrak{h}_i(qt) - \sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - \zeta_1 + \zeta_2.$$
(3.2.9)

Substituting equation (3.2.2) and equations (3.2.6) - (3.2.9), in equation (3.2.1), we have

$$\begin{split} \sum_{i=0}^{2^{J+1}} a_i \mathfrak{h}_i(t) = &\Omega\left(t, \Big[\sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(t) + \Big(-\sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - \zeta_1 + \zeta_2\Big)t + \zeta_1\Big], \\ &\Big[\sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(qt) + \Big(-\sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - \zeta_1 + \zeta_2\Big)(qt) + \zeta_1\Big], \\ &\Big[\sum_{i=0}^{2^{J+1}} a_i I_1 \mathfrak{h}_i(t) - \sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - \zeta_1 + \zeta_2\Big], \\ &\Big[\sum_{i=0}^{2^{J+1}} a_i I_1 \mathfrak{h}_i(qt) - \sum_{i=0}^{2^{J+1}} a_i I_2 \mathfrak{h}_i(1) - \zeta_1 + \zeta_2\Big]\Big). \end{split}$$
(3.2.10)

Discretisation of the above equation using collocation procedure leads to the

algebraic system for $a_i^{,s}$. After determining $a_i^{,s}$ using Newton's iterative method or any other suitable method we get the approximate solution from equation (3.2.6).

3.3 Flow Chart



3.4 Illustrative Examples

Problem 1. Consider the boundary value problem

$$y''(t) - 1 - 2(1 + t^2/8)\cos(t/2) + 2\cos(t/2)y(t/2) = 0, \ t \in [0, 1],$$
with Dirichlet boundary condition $y(0) = 1, \ y(1) = \frac{3}{2} + \sin(1).$
(3.4.1)

Whereas, the exact solution of the problem is given by $y(t) = \frac{t^2}{2} + sin(t) + 1$.

Problem 2. Consider the boundary value problem

$$y''(t) + 2e^{-t} - y(t)/2 - e^{-t/2}y(t/2) = 0, \ t \in [0, 1],$$
(3.4.2)
Dirichlet boundary condition $y(0) = 0, \ y(1) = e^{-1},$

whose exact solution is $y(t) = te^{-t}$.

with

with

Problem 3. Let us assume the boundary value problem

$$y''(t) - 4e^{-t/2}sin(t/2)y(t/2) = 0, \ t \in [0, 1],$$
(3.4.3)

along with Dirichlet boundary condition y(0) = 1, $y(1) = e^{-1}cos(1)$,

corresponding to the exact solution $y(t) = e^{-t}cos(t)$.

Problem 4. Assume the nonlinear two-point boundary value problem as

$$y''(t) - ((y(t))^{2} + (y(t))^{3})y(t/2) = 0, \ t \in [0, 1],$$
(3.4.4)
Dirichlet boundary condition $y(0) = 1, \ y(1) = 1/2.$

The exact solution of the above nonlinear boundary value problem is $y(t) = \frac{1}{t+1}$.

Problem 5. Assume the following two-point boundary value problem

$$y''(t) - y'(t)y(t/2) + 8t^2y(t/2) + \Theta(t) = 0, \ t \in [0, 1],$$
(3.4.5)

subject to Dirichlet boundary condition to y(0) = 1, y(1) = 3,

and the $\ominus(t)$ is chosen such that the exact solution of problem (3.4.5) is $y(t) = 1 + t + t^3$.

3.5 Results and Discussions

The above problems are solved using the Haar wavelet series method. For each problem at different resolutions J, the maximum absolute error and experimental convergence rate are computed. These results are presented in Tables 3.1-3.8. Table






FIGURE 3.2: Plot for Problem 2 at J=5



FIGURE 3.3: Solution curves for Problem 3 at J=5



FIGURE 3.4: loglog Plot for Problem 4 at J=5



FIGURE 3.5: loglog Plot for Problem 5 at J=5 $\,$

TABLE 3.1:	MAEs	for	Problems	1	to	5.
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J	Problem 1	Problem 2	Problem 3	Problem 4	Problem 5
1	1.5997E - 04	5.5835E - 04	5.8345E - 04	9.4871E - 04	2.6907E - 03
2	4.4574E - 05	1.5644E - 04	1.6306E - 04	2.7026E - 04	6.6283E - 04
3	1.1324E - 05	3.9757E - 05	4.1522E - 05	7.0508E - 05	1.6553E - 04
4	2.8500E - 06	1.0014E - 05	1.0454E - 05	1.7787E - 05	4.1311E - 05
5	7.1382E - 07	2.5082E - 06	2.6172E - 06	4.4582E - 06	1.0331E - 05
6	1.7831E - 07	6.2735E - 07	6.5449E - 07	1.1152E - 06	2.5828E - 06
7	4.4275E - 08	1.5692E - 07	1.6362E - 07	2.7884E - 07	6.4569E - 07
8	1.1342E - 08	3.9226E - 08	4.0890E - 08	6.9712E - 08	1.6142E - 07
9	2.6235E - 09	9.8905E - 09	1.0218E - 08	1.7427E - 08	4.0355E - 08
10	6.9708E - 10	2.4508E - 09	2.5589E - 09	4.3559E - 09	1.0088E - 08

TABLE 3.2: Rate of convergence $R_c = \frac{log\left(\frac{error(J-1)}{error(J)}\right)}{log(2)}$ for Problems 1 to 5.

J	Problem 1	Problem 2	Problem 3	Problem 4	Problem 5
1					
2	1.8435	1.8356	1.8392	1.8116	2.0213
3	1.9768	1.9763	1.9735	1.9385	2.0015
4	1.9903	1.9892	1.9898	1.9870	2.0025
5	1.9973	1.9973	1.9980	1.9963	1.9995
6	2.0012	1.9993	1.9996	1.9992	2.0000
7	2.0098	1.9992	2.0000	1.9998	2.0000
8	1.9648	2.0001	2.0005	2.0000	2.0000
9	2.1121	1.9877	2.0006	2.0001	2.0000
10	1.9121	2.0128	1.9975	2.0003	2.0001

Problem	HWS	$\operatorname{Bica}[17]$
1	1.7831E - 07 at $J = 6$	$6.9097E - 07$ at $h = \frac{\pi}{400}$
1	6.9708E - 10 at $J = 10$	$6.9100E - 09$ at $h = \frac{\pi}{4000}$
2	9.8905E - 09 at $J = 9$	$1.2770E - 08$ at $h = \frac{1}{1000}$

TABLE 3.3: Comparison of errors.

TABLE 3.4: Comparison between approximate and analytic solution at J=6.

	Problem 1		
t	y_{exact}	y_{approx}	$ y_{exact} - y_{approx} $
0.1	1.1108	1.1108	4.9295E - 08
0.2	1.2270	1.2270	9.3465E - 08
0.3	1.3415	1.3415	1.2867E - 07
0.4	1.4725	1.4725	1.5786E - 07
0.5	1.6098	1.6098	1.7498E - 07
0.6	1.7524	1.7524	1.7727E - 07
0.7	1.8995	1.8995	1.6214E - 07
0.8	2.0385	2.0385	1.3054E - 07
0.9	2.1919	2.1919	7.5335E - 08

TABLE 3.5: Comparison between approximate and analytic solution at J=6.

	Problem 2		
t	y_{exact}	y_{approx}	$ y_{exact} - y_{approx} $
0.1	9.4912E - 02	9.4912E - 02	2.8670E - 07
0.2	1.6832E - 01	1.6832E - 01	4.7186E - 07
0.3	2.2265E - 01	2.2265E - 01	5.7505E - 07
0.4	2.6907E - 01	2.6907E - 01	6.2395E - 07
0.5	3.0444E - 01	3.0444E - 01	6.1674E - 07
0.6	3.3048E - 01	3.3048E - 01	5.6142E - 07
0.7	3.4864E - 01	3.4864E - 01	4.6471E - 07
0.8	3.5953E - 01	3.5953E - 01	3.4350E - 07
0.9	3.6601E - 01	3.6601E - 01	1.8183E - 07

	Problem 3		
t	y_{exact}	y_{approx}	$ y_{exact} - y_{approx} $
0.1	8.9490E - 01	8.9490E - 01	3.0232E - 07
0.2	7.9563E - 01	7.9563E - 01	4.9807E - 07
0.3	7.0701E - 01	7.0701E - 01	6.0542E - 07
0.4	6.1535E - 01	6.1535E - 01	6.5268E - 07
0.5	5.2907E - 01	5.2907E - 01	6.3862E - 07
0.6	4.4879E - 01	4.4879E - 01	5.7344E - 07
0.7	3.7491E - 01	3.7491E - 01	4.6661E - 07
0.8	3.1255E - 01	3.1255E - 01	3.3853E - 07
0.9	2.5139E - 01	2.5139E - 01	1.7510E - 07

TABLE 3.6: Comparison between approximate and analytic solution at J=6.

TABLE 3.7: Comparison between approximate and analytic solution at J=6.

	Problem 4		
t	y_{exact}	y_{approx}	$ y_{exact} - y_{approx} $
0.1	9.0459E - 01	9.0459E - 01	7.0007E - 07
0.2	8.2848E - 01	8.2848E - 01	1.0092E - 06
0.3	7.6877E - 01	7.6877E - 01	1.1088E - 06
0.4	7.1309E - 01	7.1309E - 01	1.0937E - 06
0.5	6.6494E - 01	6.6493E - 01	9.9756E - 07
0.6	6.2287E - 01	6.2287E - 01	8.4864E - 07
0.7	5.8581E - 01	5.8581E - 01	6.6355E - 07
0.8	5.5531E - 01	5.5531E - 01	4.6913E - 07
0.9	5.2567E - 01	5.2567E - 01	2.3844E - 07

TABLE 3.8: Comparison between approximate and analytic solution at J=6.

	Problem 5		
t	y_{exact}	y_{approx}	$ y_{exact} - y_{approx} $
0.1	1.1066	1.1066	7.4140E - 07
0.2	1.2159	1.2159	1.3630E - 06
0.3	1.3280	1.3280	1.8446E - 06
0.4	1.4675	1.4675	2.2490E - 06
0.5	1.6319	1.6319	2.5062E - 06
0.6	1.8274	1.8274	2.5806E - 06
0.7	2.0605	2.0605	2.4239E - 06
0.8	2.3143	2.3143	2.0191E - 06
0.9	2.6371	2.6371	1.2202E - 06

3.1 demonstrate that error decreases with increase in resolution(J). In Table 3.2, we have calculated the rate of convergence by using the formula

$$R_c = \frac{\log(\frac{E_{J-1}}{E_J})}{\log(2)}$$

where E_J is the maximum absolute errors at resolution J. Also, from Table 3.2 one can observe that the experimental convergence rates in each problem tend to be 2, as described in [67]. The approximate and analytical solution curves of each problem are plotted in figures 3.1-3.5. We observed that both curves coincide, and a comparison drawn in Table 3.3 shows the method is more accurate and efficient.

3.6 Summary

This chapter¹ introduced the Haar wavelet-based numerical method to solve Dirichlet boundary value problems of proportional delay nature. The method is tested on benchmark problems and the results are compared with the existing method [17]. The accuracy and convergence rate of the technique have been validated by solving both linear and nonlinear problems. The MATLAB package is utilized to perform computer simulations, and the solutions obtained were compared with the analytical solutions.

¹The content of this chapter is published in **International Journal of Nonlinear Analysis** and **Applications.**(ESCI, Scopus Indexed).

Chapter 4

Haar wavelet series method for solving simultaneous proportional delay differential equations

4.1 Introduction

In this study, we apply the Haar wavelet series method (HWSM) to solve the following simultaneous proportional delay differential equations:

$$y_1'(t) = \Omega_1(t, y_1(t), y_2(t), \dots, y_n(t), y_1(q_1t), y_2(q_2t), \dots, y_n(q_\lambda t))$$

$$y_2'(t) = \Omega_2(t, y_1(t), y_2(t), \dots, y_n(t), y_1(q_1t), y_2(q_2t), \dots, y_n(q_\lambda t))$$

$$\vdots \qquad (4.1.1)$$

$$y'_{n}(t) = \Omega_{n}(t, y_{1}(t), y_{2}(t), \dots, y_{n}(t), y_{1}(q_{1}t), y_{2}(q_{2}t), \dots, y_{n}(q_{\lambda}t))$$

$$y_{\sigma}(0) = y_{\sigma 0}, \quad \sigma = 1, 2, \dots, n,$$

where Ω_{σ} 's are analytical functions, and q_{σ} 's $\in (0, 1), \sigma = 1, 2, ..., n$. These equations have been widely noticed in several models, such as biological models, aerospace systems, control theory, disease spread models, tumour growth models, etc. Therefore, this study is an important contribution in the field of applied mathematical modelling and numerical analysis. This method utilizes delayed Haar wavelet series and collocation points to transform the simultaneous proportional delay differential equations into a system of algebraic matrix equations with unknown coefficient matrices. The values of these unknown row matrices can be obtained by using a suitable solver. With these coefficients, the solution in terms of the collocated Haar wavelet series is obtained.

4.2 Description of Method

In this section, we apply the Haar wavelet series method for solving simultaneous proportional delay differential equation (4.1.1).

Let

$$y'_{1}(t) = \sum_{i=1}^{2^{J+1}} a_{i} \mathfrak{h}_{i}(t),$$

$$y'_{2}(t) = \sum_{i=1}^{2^{J+1}} b_{i} \mathfrak{h}_{i}(t),$$

$$\vdots$$
(4.2.1)

$$y'_n(t) = \sum_{i=1}^{2^{J+1}} c_i \mathfrak{h}_i(t),$$

where a_i 's, b_i 's... c_i 's are Haar wavelet coefficients.

Now, by integrating equation (4.2.1), we have

$$y_{1}(t) = \sum_{i=1}^{2^{J+1}} a_{i} I_{1} \mathfrak{h}_{i}(t) + y_{1}(0),$$

$$y_{2}(t) = \sum_{i=1}^{2^{J+1}} b_{i} I_{1} \mathfrak{h}_{i}(t) + y_{2}(0),$$
(4.2.2)

$$y_n(t) = \sum_{i=1}^{2^{J+1}} c_i I_1 \mathfrak{h}_i(t) + y_n(0).$$

÷

Also,

$$y_{1}(q_{1}t) = \sum_{i=1}^{2^{J+1}} a_{i}I_{1}\mathfrak{h}_{i}(q_{1}t) + y_{1}(0),$$

$$y_{2}(q_{2}t) = \sum_{i=1}^{2^{J+1}} b_{i}I_{1}\mathfrak{h}_{i}(q_{2}t) + y_{2}(0),$$

$$\vdots$$

$$(4.2.3)$$

$$y_n(q_n t) = \sum_{i=1}^{2^{n+1}} c_i I_1 \mathfrak{h}_i(q_n t) + y_n(0),$$

where $y_{\sigma}(0)$'s are given initial conditions and q_{σ} 's $\in (0, 1), \sigma = 1, 2, ..., n$. Now, upon plugging equations (4.2.1), (4.2.2) and (4.2.3) in the governing equation (4.1.1) along with the collocation points we get a system of an algebraic equations. After solving the system we obtain the unknowns Haar wavelet coefficients and finally using these coefficients in equation (4.2.2) approximate solution at collocation points can be obtained.

4.3 Algorithm

Input: Level of resolution J.

Step-1: Set collocation points $t_l = \frac{(l-0.5)}{2^{J+1}}, \quad l = 1, 2, 3 \dots 2^{J+1}.$

Step-2: Compute Haar wavelets \mathfrak{h}_i and integral of Haar wavelets $I_1\mathfrak{h}_i$ from equations

(2.2.3) and (2.2.9), respectively.

Step-3: Assume $y'_{\sigma}(t) = \sum_{i=1}^{2^{J+1}} (a_i)_{\sigma} \mathfrak{h}_i(t)$ where $\sigma = 1, 2, \dots, n$ and i is wavelet index.

Step-4: Integration of Step-3 within the limits 0 to t yields,

$$y_{\sigma}(t) = \sum_{i=1}^{2^{J+1}} (a_i)_{\sigma} I_1 \mathfrak{h}_i(t) + y_{\sigma}(0), \quad \sigma = 1, 2, \dots, n.$$

Step-5: Upon using the expressions for $y'_{\sigma}(t)$'s, $y_{\sigma}(t)$'s ,..., $y_{\sigma}(q_{\sigma}t)$'s along with collocation points in equation (4.1.1), we get an algebraic system in unknown vectors $(a_i)_1, (a_i)_2, \ldots, (a_i)_n$.

Step-6: Solve the algebraic system for the unknown vectors $(a_i)_1, (a_i)_2, \ldots, (a_i)_n$.

Step-7: Utilize the values of the vectors $(a_i)_1, (a_i)_2, \ldots, (a_i)_n$ in Step-4 to obtained the unknown functions $y_1(t), y_2(t), \ldots, y_n(t)$.

Output: Approximate solution $y_1(t_l), y_2(t_l), \ldots, y_n(t_l)$ is obtained.

4.4 Implementation of Method on Test Problems

The efficiency and numerical validation of the method is demonstrated with the help of following test problems.

Problem 1: Consider the system

$$\begin{cases} y_1'(t) = y_1(t/2) + y_1(t) - y_2(t) + exp(-t) - exp(t/2) \\ y_2'(t) = -y_1(t/2) - y_1(t) - y_2(t) + exp(t) + exp(t/2) \\ with \ y_1(0) = y_2(0) = 1, \quad 0 \le t \le 1. \end{cases}$$

$$(4.4.1)$$

The system has exact solution $y_1(t) = \exp(t), y_2(t) = \exp(-t)$.

We have solved the above system of differential equation using the Haar wavelet series method. The maximum absolute error (MAE) at different levels of resolution is given in the Table 4.1. The Table 4.1 shows that the error decreases from O(10-2) to O(10-7), ensuring the method's convergence. Also, Fig 4.1 shows that approximate and exact curves match closely.

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FIGURE 4.1: Exact and approximate solution of Problem 1.

	Problem 1		Problem 2	
J	$\mathbf{MAE:} y_1(t)$	$\mathbf{MAE:} y_2(t)$	$\mathbf{MAE:} y_1(t)$	$\mathbf{MAE:} y_2(t)$
2.0	1.3000E - 02	5.4000E - 03	7.3000E - 03	7.3000E - 03
3.0	3.4187E - 03	1.4267E - 03	1.9007E - 03	1.9007E - 03
4.0	8.7870E - 04	3.6790E - 04	4.8190E - 04	4.8190E - 04
5.0	2.2290E - 04	9.3500E - 05	1.2170E - 04	1.2170E - 04
6.0	5.6150E - 05	2.3570E - 05	3.0580E - 05	3.0570E - 05
7.0	1.4090E - 05	5.9200E - 06	7.6620E - 06	7.6620E - 06
8.0	3.5290E - 06	1.4820E - 06	1.9180E - 06	1.9180E - 06
9.0	8.8310E - 07	3.7090E - 07	4.7980E - 07	4.7980E - 07

TABLE 4.1: MAEs at different levels of J.

Problem 2: Consider the system

$$\begin{cases} y_1'(t) = \exp(t/2)y_2(t/2) + y_1(t) \\ y_2'(t) = \exp(t/2)y_1(t/2) + y_2(t) \\ with \ y_1(0) = y_2(0) = 1, \ 0 \le t \le 1. \end{cases}$$

$$(4.4.2)$$

The system has exact solution $y_1(t) = \exp(t), y_2(t) = \exp(t)$.





FIGURE 4.2: Exact and approximate solution of Problem 2.



FIGURE 4.3: Absolute error vs J.

The approximate solution of the aforementioned problem obtained using the Haar wavelet series approach is satisfactory. Computer simulation is carried out and the numerical results are presented in Table 4.1. The exact and approximate solution of Problem 2 is plotted in Fig 4.2 for J=4. Also, the convergence of the method can be observed in Fig 4.3.

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mate solution of Problem 3.

FIGURE 4.5: Absolute error vs J (Problem 3).

Problem 3: Consider the following non-linear proportional delay system:

$$\begin{cases} y_1'(t) = 2y_2(t/2) + y_3(t) + t\cos(t/2) \\ y_2'(t) = -2y_3^2(t) + 1 - t\sin(t) \\ y_3'(t) = y_2(t) - y_1(t) - t\cos(t) \\ with \ y_1(0) = -1, y_2(0) = y_3(0) = 0, \quad 0 \le t \le 1. \end{cases}$$

$$(4.4.3)$$

Exact solution is $y_1(t) = -\cos(t), y_2(t) = t\cos(t), y_1, (t) = \sin(t)$. In this case we have solved a non-linear system with three equations using the Haar wavelet series method. The exact and approximate solution for Problem 3 is plotted in Fig. 4.4. It is observed that both curves match closely. Absolute errors at different levels of resolution are depicted in Table 4.2 which shows that the computed numerical solution is satisfactory. Also, it is evident from Fig.4.5 that the accuracy of the solution is directly proportional to the value of J. Furthermore, in each case, the error was reduced from O(10-2) to O(10-7).

J	$\mathbf{MAE:} y_1(t)$	$\mathbf{MAE:}y_2(t)$	$\mathbf{MAE:}y_{3}(t)$
2.0	1.9167E - 03	2.0504E - 03	3.1668E - 03
3.0	4.8600E - 04	5.1360E - 04	8.0760E - 04
4.0	1.2190E - 04	1.2850E - 04	2.0340E - 04
5.0	3.0510E - 05	3.2120E - 05	5.1000E - 05
6.0	7.6300E - 06	8.0300E - 06	1.2770E - 05
7.0	1.9070E - 06	2.0070E - 06	3.1930E - 06
8.0	4.7680E - 07	5.0190E - 07	7.9860E - 07
9.0	1.1920E - 07	1.2550E - 07	1.9970E - 07

TABLE 4.2: MAEs at different levels of J (Problem 3).

4.5 Summary

This chapter¹ discussed the Haar wavelet series method for linear and nonlinear simultaneous proportional delay differential equations with two and three dependent variables. Three illustrations at different counts of collocation points are solved to test the consistency, correctness, and efficacy of the method. The numerical results demonstrated in figures and tables illustrate the expected rate of convergence.

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Chapter 5

A modified Haar wavelet series method to solve higher-order multi-pantograph equations arising in electrodynamics

5.1 Pantograph Equations

The functional-differential with proportional delay is known as the pantograph equation or generalized pantograph equations. The name pantograph first appeared in 1851 and was a device used in the construction of the electric locomotive. The mathematical model of the pantograph was first developed by Ockendon and Tyler [79]. The Pantograph equation is one of the most distinguished delay differential equations and has been an interest of many researchers [19, 35, 36]. The pantograph differential equations are encountered in studies of population dynamic models, quantum theory, control theory, cell growth models, disease spread models, and astrophysics [35]. These equations also have several industrial applications and play a central role in the mathematical modelling of the train's overhead current collection system [79]. The continuous electricity supply between the catenary and the train's motor is maintained by a device called a z-shape pantograph. The z-shape pantograph (also known as half pantograph) resembled the pantograph device for copying, writing, and drawing. It has a spring mechanism that pushes the contact shoe up against the wire to draw the electricity required to power the train.

Most of these equations can not be solvable exactly. Therefore, a numerical technique is required to obtain their approximate solutions. Variational iteration method [25], One leg- θ method [106], Two-stage R-K method [108], Reproducing kernel Hilbert space method (RKHSM) [64], Differential transform method [53], Adomain decomposition method [20], Perturbed iteration method [10] are some well known numerical techniques to solve such types of differential equations. Recently, in [8] time-invariant and time-varying first-order delay differential equations have been solved using the Haar wavelet collocation method. Some other collocation methods are also developed using Chebyshev Polynomials, Hermite Polynomials, and Bernoulli Polynomials. For details, readers may refer to [63, 103, 113] respectively.

Chen and Hsiao [24] gave an idea of utilizing Haar operational matrix of integration for solving differential equations. In the existing literature, the development and application of the Haar wavelet collocation method (HWCM) for solving differential equations are based on the method given by Chen and Hsiao. Later this idea has been extended to solve a wide range of problems [60]. Marzban and Razzaghi [71] adapted the rationalized Haar wavelet approach for solving nonlinear optimal control problems. The HWCM is also a efficient tool in structural mechanics, Hariharan [46] applied the Haar wavelet-based technique for solving the finite length beam equation. Lepik [60] discussed buckling of elastic beams using the Haar wavelet method. In [84] Patra and Saha obtained the solution of stiff point kinetics equations using wavelet operational method based on the Haar wavelet. In the recent past, the Chen-Hsiao technique has been extended to solve delay differential equations. Aziz and Amin^[8] investigated the approximate solution of delay differential as well as partial delay differential equations. Raza et al. ^[89] transformed the delay term using Taylor series expansion and then applied the Haar wavelet collocation method to solve singularly perturbed differential-difference equations and singularly perturbed convection delayed dominated diffusion equations. Abdullah and Rafiq ^[2] combined the backward Euler method and the HWCM to obtain the approximate solution of the Chen-Lee-Liu equation.

Here we have applied a modified Haar wavelet series method (MHWSM) instead of the conventional Haar wavelet collocation method. Instead of the highest (n^{th}) order derivative, we expanded the $(n + 1)^{th}$ order derivative involved in the differential equations in terms of the Haar series. The MHWSM produced a smoother solution than the Haar wavelet collocation method, therefore a significant decrease in absolute error is expected.

5.2 Construction of Method

Let us assume a n^{th} order pantograph equation of the form

$$y^{n}(t) = \varphi(g(t), y(q_{0}t), y^{1}(q_{1}t), y^{2}(q_{2}t) \dots y^{n}(q_{n}t)), \quad \forall \ t \in [t_{0}, t_{f}]$$
with $y^{\eta}(t_{0}) = y_{0}^{\eta},$
(5.2.1)

where $\varphi : [t_0, t_f] \times \underbrace{R \times R \cdots \times R}_{(n+1)\text{-times}} \mapsto R$ is a differentiable function, g(t) is continuous

on $[t_0, t_f]$ and $q_0, q_1 \dots q_n$ are real constants lies in (0, 1]. Also, $y^1, y^2 \dots y^n$ denotes the first, second and n^{th} order derivatives, respectively and y_0^{η} are initial value conditions $\eta = 0, 1, 2 \dots n - 1$. Put $t = t_0$ in equation (5.2.1) for $y^n(t_0)$. In order to solve equation (5.2.1) we have established the following algorithm using Haar wavelet series equation (2.2.3). Let us suppose $y^{n+1}(t)$ be square integrable function. Therefore, we can write

$$y^{n+1}(t) = \sum_{i=1}^{2M} a_i \mathfrak{h}_i(t).$$
 (5.2.2)

Integrating equation (5.2.2) r times with respect to t, we have the following relation

$$y^{n+1-r}(t) = \sum_{i=1}^{2M} a_i I_r \mathfrak{h}_i(t) + \sum_{\eta=n+1-r}^n \frac{y^{\eta}(t_0)(t)^{\eta-(n+1-r)}}{(\eta-(n+1-r))!}.$$
 (5.2.3)

Taking r = n + 1 in relation equation (5.2.3), we have y(t) as

$$y(t) = \sum_{i=1}^{2M} a_i I_{n+1} \mathfrak{h}_i(t) + \sum_{\eta=0}^n \frac{y^{\eta}(t_0)(t)^{\eta}}{(\eta)!}.$$
 (5.2.4)

Also,

$$y(q_0 t) = \sum_{i=1}^{2M} a_i I_{n+1} \mathfrak{h}_i(q_0 t) + \sum_{\eta=0}^n \frac{y^{\eta}(t_0)(q_0 t)^{\eta}}{(\eta)!}.$$
 (5.2.5)

Similarly,

$$y^{1}(q_{1}t) = \sum_{i=1}^{2M} a_{i}I_{n}\mathfrak{h}_{i}(q_{1}t) + \sum_{\eta=1}^{n} \frac{y^{\eta}(t_{0})(q_{1}t)^{\eta-1}}{(\eta-1)!}, \qquad (5.2.6)$$

:

$$y^{n}(q_{n}t) = \sum_{i=1}^{2M} a_{i}I_{1}\mathfrak{h}_{i}(q_{n}t) + y^{n}(t_{0}). \qquad (5.2.7)$$

Now, substituting equations (5.2.3 - 5.2.7) in equation (5.2.1), we get

$$\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(t) + y^n(t_0) = \varphi \left(g(t), \sum_{i=1}^{2M} a_i I_{n+1} \mathfrak{h}_i(q_0 t) + \sum_{\eta=0}^n \frac{y^\eta(t_0)(q_0 t)^\eta}{(\eta)!}, \\ \sum_{i=1}^{2M} a_i I_n \mathfrak{h}_i(q_1 t) + \sum_{\eta=1}^n \frac{y^\eta(0)(q_1 t)^{\eta-1}}{(\eta-1)!}, \dots, \sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(q_n t) + y^n(0) \right).$$
(5.2.8)

Moreover, utilizing collocation points $t_l = \frac{l-0.5}{2M}, l = 1, 2, \dots 2M$ in equation (5.2.8), we obtain

$$\sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(t_l) + y^n(t_0) = \varphi \left[g(t_l), \sum_{i=1}^{2M} a_i I_{n+1} \mathfrak{h}_i(q_0 t_l) + \sum_{\eta=0}^n \frac{y^\eta(t_0)(q_0 t_l)^\eta}{(\eta)!}, \\ \sum_{i=1}^{2M} a_i I_n \mathfrak{h}_i(q_1 t_l) + \sum_{\eta=1}^n \frac{y^\eta(t_0)(q_1 t_l)^{\eta-1}}{(\eta-1)!}, \dots, \sum_{i=1}^{2M} a_i I_1 \mathfrak{h}_i(q_n) + y^n(t_0) \right].$$
(5.2.9)

Now, one can determine the coefficients $a_i^{i}s$ by using any iterative techniques like Newton's method or Broyden's method. Finally, the solution is obtained by substituting $a_i^{i}s$ in equation (5.2.4).

5.3 Algorithm

Input: level of resolution J.

- Step-1: Set collocation points $t_l = \frac{(l-0.5)}{2M}$, $l = 1, 2, 3 \dots 2M, M = 2^J$.
- Step-2: Compute Haar wavelets $\mathfrak{h}_i(t)$ and integral of Haar wavelets $I_n\mathfrak{h}_i(t)$ from equations (1.2.1) and (1.2.2), respectively.
- Step-3: Construct the system (5.2.9) by using Eqs. (5.2.3-5.2.7) and collocation points t_l sets in step 1.

Step-4: Apply Newton's method to the system (5.2.9) for unknowns a_i 's.

Step-5: Put $a_i s$ in Eq. (5.2.4).

Output: Approximate solution $y_h(t_l)$.

5.4 Numerical Experiment

To check the applicability and efficiency of our technique, we have solved secondorder linear and non-linear differential equations, integro differential equations, a third-order, and a fourth-order differential equation of pantograph nature. All computer simulations are carried out in MATLAB and are reported in tables and figures.

Problem 1. Consider the pantograph equation

$$y''(t) = \frac{3}{4}y(t) + y\left(\frac{t}{2}\right) + y'\left(\frac{t}{2}\right) + \frac{1}{2}y''\left(\frac{t}{2}\right) - t^2 - t + 1, \ t \in [0, 1],$$

(5.4.1)
$$y(0) = 0, \ y'(0) = 0.$$

The exact solution of equation (5.4.1) from [64] is $y_e = t^2$.

The present technique is successfully applied on equation (5.4.1) and the result is compared with some existing methods [25, 64, 106, 107, 108]. Wavelet coefficients are calculated using the classical Newton's method by choosing an appropriate initial guess. We observed that the maximum absolute error is zero for J =2, 3, 4, ..., 9. Also, it is evident from Table 5.1 and Table 5.2 that our method has easy applicability and produces better results. Figure 5.1 shows that both exact and approximate solutions coincide visually.



FIGURE 5.1: Comparison of exact and approximate solution (Problem 1).

TABLE 5.1 :	Error	comparison	with	existing	method	(Problem	1)
---------------	-------	------------	------	----------	--------	----------	----

Our	Two-stage[108]	One-leg	VIM[25]	VIM[25]	RKHSM
\mathbf{method}	order-one	$\theta ext{-method}$	n = 5	n = 5	n = 100[64]
J=2	RKM	$(\theta = 0.8)[106]$			
0	5.34E - 03	2.81E - 01	1.11E - 02	5.55E - 03	4.92E - 04

Problem 2. In this Example we consider a second order nonlinear pantograph equation

$$y''(t) = -y(t) + 5\left(y\left(\frac{t}{2}\right)\right)^2, \ t \in [0, 1],$$

$$y(0) = 1, \ y'(0) = -2.$$

(5.4.2)

Approximate solution of equation (5.4.2) is obtained with the present algorithm. Our solution is compared with exact solution $y_e = \exp(-2t)$ in Table 5.3 and Figure 5.2. We have observed that maximum absolute errors are decreased from order of 10^{-3} for J = 2 to order of 10^{-7} for J = 9.

t(=1/32)	Present method	Exact solution
1	0.00097656250	0.00097656250
3	0.00878906250	0.00878906250
5	0.02441406250	0.02441406250
7	0.04785156250	0.04785156250
9	0.07910156250	0.07910156250
11	0.11816406250	0.11816406250
13	0.16503906250	0.16503906250
15	0.21972656250	0.21972656250
17	0.28222656250	0.28222656250
19	0.35253906250	0.35253906250
21	0.43066406250	0.43066406250
23	0.51660156250	0.51660156250
25	0.61035156250	0.61035156250
27	0.71191406250	0.71191406250
29	0.82128906250	0.82128906250
31	0.93847656250	0.93847656250

TABLE 5.2: Comparison of exact and approximate solution (Problem 1).

TABLE 5.3 :	MAEs at different levels
of	J (Problem 2).

TABLE 5.4: MAEs at different levels of J (Problem 4).

J	$\max y_{exact} - y_{approx} $	J	$\max y_{exact} - y_{approx} $
3	7.4217E - 04	3	8.2836E - 04
4	1.9187E - 04	4	2.1882E - 04
5	4.8675E - 05	5	5.6115E - 05
6	1.2252E - 05	6	1.4201E - 05
7	3.0729E - 06	7	3.5715E - 06
8	7.6943E - 07	8	8.9555E - 07
9	1.9248E - 07	9	2.2414E - 07



FIGURE 5.2: Comparison of exact and approximate solution (Problem 2).



FIGURE 5.3: Maximum absolute errors vs J (Problem 2).

Problem 3. Let us consider the following second order pantograph type initial value problem mention in [64],

$$y''(t) = y'\left(\frac{t}{2}\right) - \frac{t}{2}y''\left(\frac{t}{2}\right) + 2, \ t \in [0, 1],$$

$$y(0) = 1, \ y'(0) = 0.$$

(5.4.3)

We have solve this example using present method. The approximate function to be sought is $y_h = \sum_{i=1}^{2M} a_i * I_3 \mathfrak{h}_i(t) + t^2 + 1$. A computer simulation is carried out and it is observed that the maximum absolute error is zero for $J = 2, 3, \ldots, 9$. A comparison

t(=1/16)	Present method	Exact solution	
1	1.003906250	1.003906250	
3	1.035156250	1.035156250	
5	1.097656250	1.097656250	
7	1.191406250	1.191406250	
9	1.316406250	1.316406250	
11	1.472656250	1.472656250	
13	1.660156250	1.660156250	
15	1.878906250	1.878906250	

TABLE 5.5: Comparison of exact and approximate solution (Problem 3).





between the approximate solution and the exact solution is demonstrated in Table 5.5 and Figure 5.4, which show that both solutions coincide.

Problem 4. We consider a nonlinear integro-differential equation with proportional delay in kernal

$$y'(t) + \left(\frac{t}{2} - 2\right) y(t) - 2 \int_0^t \left(y\left(\frac{s}{2}\right)\right)^2 ds = 1, \ t \in [0, 1],$$

$$y(0) = 0.$$
 (5.4.4)

t(=1/32)	Present method	Exact solution	$ y_{exact} - y_{approx} $
1	0.0322421468	0.0322419814	0.1653E - 6
3	0.1029618079	0.1029642319	0.2423E - 5
5	0.1826661951	0.1826747572	0.8562E - 5
7	0.2722202080	0.2722387735	0.1856E - 4
9	0.3725622277	0.3725957133	0.3348E - 4
11	0.4847113430	0.4847651995	0.5385E - 4
13	0.6097726143	0.6098534812	0.8086E - 4
15	0.7489450753	0.7490603671	0.1152E - 3
17	0.9035281845	0.9036866916	0.1585E - 3
19	1.0749307939	1.0751423553	0.2115E - 3
21	1.2646789302	1.2649549829	0.2760E - 3
23	1.4744259044	1.4747792434	0.3533E - 3
25	1.7059615777	1.7064068834	0.4453E - 3
27	1.9612238562	1.9617775258	0.5536E - 3
29	2.2423096473	2.2429902915	0.6806E - 3
31	2.5514879480	2.5523163047	0.8283E - 3

TABLE 5.6: Comparison of exact and approximate solution (Problem 4).

Equation (5.4.4) can be reduced to following second order nonlinear pantograph equation

$$y''(t) + \left(\frac{t}{2} - 2\right)y'(t) + \frac{1}{2}y(t) - 2\left(y\left(\frac{t}{2}\right)\right)^2 = 1, \ t \in [0, 1],$$

$$y(0) = 0, \ y'(0) = 1.$$

(5.4.5)

Now we have applied the present algorithm to equation (5.4.5) and obtained its approximate solution. The exact solution $y_e = t \exp(t)$, is compared to the approximate solution and results are shown in Table 5.4 and Table 5.6. We have observed that maximum absolute errors are decreased from order of 10^{-4} for J = 3to order of 10^{-7} for J = 9. We have verified in Figure 5.5 that both solutions visually coincide.



FIGURE 5.5: Comparison of exact and approximate solution (Problem 4).



FIGURE 5.6: Absolute errors vs J (Problem 4).

TABLE 5.7: Error comparison with existing method (Problem 5).

Our method	Two-stage[108]	VIM[25]	VIM[25]	VIM[25]
J=2	order-one RKM	n = 4	n = 5	n = 6
6.92214E - 10	7.34E - 02	3.21E - 04	4.01E - 05	1.26E - 06

Problem 5. In this Example we consider a third-order pantograph equation

$$y'''(t) = y(t) + y'\left(\frac{t}{2}\right) + y''\left(\frac{t}{3}\right) + \frac{1}{2}y'''\left(\frac{t}{4}\right) - t^4 - \frac{t^3}{2} - \frac{4}{3}t^2 + 21t, \ t \in [0, 1],$$

$$y(0) = y'(0) = y''(0) = 0.$$

(5.4.6)

TABLE 5.8: Comparison of exact and approximate solution (Problem 5).				
$t(=\frac{1}{16})$	Present method	Exact solution	$ y_{exact} - y_{approx} $	
1	0.000015258789160	0.000015258789063	0.00098E - 10	
3	0.001235961921205	0.001235961914063	0.07143E - 10	
5	0.009536743166434	0.009536743164063	0.02372E - 10	
7	0.036636352518005	0.036636352539063	0.21057E - 10	
9	0.100112915165484	0.100112915039063	1.26421E - 10	
11	0.223403930795185	0.223403930664063	1.31123E - 10	
13	0.435806274199234	0.435806274414063	2.14829E - 10	
15	0.772476195596849	0.772476196289063	6.92214E - 10	





FIGURE 5.7: Comparison of exact and approximate solution (Problem 5).

FIGURE 5.8: Maximum absolute errors vs J (Problem 5).

We have applied the present algorithm to equation (5.4.6). A comparison between approximate and exact solution $y_e = t^4$ demonstrated in Figure 5.7 and it shows that both solutions visually coincide. The wavelet coefficients are calculated using the classical Newton's method with an appropriate initial guess. In addition, we discovered that the maximum absolute error for J = 2 is $O(10)^{-10}$. From Tables 5.7, 5.8 and 5.9 we conclude that the present method is more efficient and produces much better results.

Problem 6. Now we consider a fourth order nonlinear multi-pantograph equation

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1		0	
TABLE 5.9: MAEs at different levels of J (Problem 5).		TABLE 5.10: MAEs at different lev- els of J (Problem 6).	
J	$\max y_{exact} - y_{approx} $	J	$\max y_{exact} - y_{approx} $
2	6.9221E - 10	4	4.3588E - 05
3	4.7252E - 10	5	1.9020E - 05
4	2.1423E - 11	6	4.8286E - 06
5	1.9649E - 11	7	2.2318E - 06
6	1.0246E - 12	8	5.5998E - 07
7	9.1538E - 13	9	2.6855E - 07

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and approximate solution (Problem 6).

FIGURE 5.10: Maximum absolute errors vs J (Problem 6).

$$y^{iv}(t) = y''\left(\frac{t}{2}\right) \left(y^{iv}\left(\frac{t}{4}\right) - y(t)\right) + \lambda(t), \ t \in [0, 1],$$

$$y(0) = 0, \ y'(0) = 1, \ y''(0) = 2, \ y'''(0) = 2,$$

(5.4.7)

where $\lambda(t)$ is supplied in such a way that the system has the exact solution $y_e = e^t sin(t)$. Carrying out the numerical technique mentioned in section 3, we have obtained the approximate solution of equation (5.4.7) for different values of J. Maximum absolute errors are computed at different resolutions (table 5.10). Moreover, the exact solution and approximate solution are plotted in figure 5.9 for J = 4. Based on the obtained results, it is realized that the method is efficient for tackling such problems.

5.5 Summary

This chapter¹ modified the traditional Haar wavelet series method and utilized the method to obtain the solutions of second and higher-order Pantograph equations. A short account of applications of pantograph equations in electrodynamics was also presented. The method was tested in numerical simulations and compared with other methods.

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Chapter 6

A collocation method for solving proportional delay Riccati differential equations of fractional order

6.1 Introduction

Over the last few years, the subject of fractional calculus has provided more generalized methods to describe the behaviours of several physical systems. Fractional order models have attracted the attention of researchers since differential equations involving non integer derivative demonstrate the dynamics of many systems more realistically, for instance bio-engineering[65], viscoelasticity[66], diffusion[72], chaos theory, Fractional capacitor theory[86], electromagnetism[114], electrochemistry[81], and many others [30, 104]. Also, in recent years, several authors have contributed a large literature on the analysis and applications of fractional differential equations, namely [32, 55, 85]. Furthermore, several authors suggested different definitions of fractional derivatives. The Riemann-Liouville and Liouville-Caputo define a fractional derivative, which has significant importance in the field of fractional calculus but has certain limitations due to singular local power kernel [55, 85]. Caputo and Fabrizio propose a new definition of fractional derivatives using exponential decay kernel^[22]. Recently, Atangana and Baleanu suggested another fractional derivative which is based on the concept of ML function [7]. Furthermore, some authors suggested more interesting definitions where time-dependent variable-order fractional derivative and integral $\alpha(t)$ is involved [95].

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Delay differential equations (DDEs) form a special class of differential equations in which the rate of solution depends on the present as well as some previous value of the dependent variable and/or their derivative. DDEs are widely used to model processes in physical sciences, biosciences, engineering, electrodynamics, and economics. Depending upon the nature of delay/lag τ , DDEs have various formats such as DDEs with constant delay, DDEs with time dependent delay $\tau(t)$, DDEs with state dependent delay $\tau(t, y)$, Neutral DDEs, and proportional DDEs or Pantograph equations. Herein, the subject of our interest is the following fractional order proportional delay Riccati differential equation(PDRDE):

$$\begin{cases} \chi^{\alpha}(t) = \Omega(t) + c_1(t)\chi(t) + c_2(t)\chi(qt)(c_3(t) - \chi(qt)), \\ with \ \chi(0) = \chi_0, \quad 0 \le t \le 1, \quad 0 < \alpha < 1, \end{cases}$$
(6.1.1)

where Ω, c_1, c_2, c_3 are analytical functions, and $q \in (0, 1)$. Note that when q = 1, equation (6.1.1) become ordinary Riccati fractional differential equation.

Due to the involvement of fractional derivatives and delay terms, the computational complexity of these equations increases, making it too complicated to solve these equations analytically. Also, in some cases, the analytical solution of these equations does not exist, so an efficient numerical technique is necessary to find the approximate solution.

In recent years, many authors have shown an interest in solving delay differential equations. Some of them are as follows. In [116] Yuzbasi & Sezer have introduced approximation method based on exponential polynomial and collocation point for solving pantograph equations, Bahsi & Cevik [11] solved proportional DDEs numerically using perturbation-iteration method, Davaeifar & Rashidinia^[29] has proposed a collocation method using first Boubaker polynomials for solving multipantograph equations. Further, in [94] Sakar et al. employed iterative reproducing kernel method for solving Riccati differential equations, Muthukumar & Ganesh [77] used shifted jacobi polynomial to develop numerical technique for solving fractional delay differential equations, Li & Wang [62] represent solution of linear fractional DDEs of Riemann-Liouville type using the Mittage-Leffler function. Ali et al. [6] used the spectral collocation method to solve fractional order DDEs, Ghomanjani & Shateyi [41] developed approximation scheme using Genocchi polynomial for solving quadratic Riccati DEs, Multipantograph DDEs and optimal control systems with pantograph delays, Jafari et al. [51] provide transferred Legendre pseudospectral method to solve pantograph DDEs, Izadi & Srivastava^[50] produced numerical solution of the Lane Emeden pantograph DDEs using Bessels polynomials and collocation points. Panghal & Kumar [83] used a neural network technique to solve multipantograph type DDEs.

The Haar wavelet is a compact support box function which takes only three values $\{0, \pm 1\}$. In the last two decades, it has been used to solve a wide variety of differential equations. The detailed applications of the Haar wavelet can be found in the monograph and references therein in [60, 87]. In [97] Shah et al. proposed a numerical technique using Haar wavelet for solving fractional differential equations. Recently, Oruc et al. [82] presented Haar wavelet and finite difference based scheme to solve the two-dimensional time fractional reaction—sub-diffusion equation. Akmal

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& Arshad [88] solved neutral DDEs using Haar wavelet bases, Abdullah & Rafiq[2] combined the backward Euler method and Haar wavelet collocation method to obtain the approximate solution of the Chen-Lee-Liu equation. Motivated by the above literature, we aim to apply the Haar wavelet series method (HWSM) to solve the proportional delay Riccati differential equation of fractional order with Caputo derivative.

Some fundamental definitions related to fractional calculus,

Definition 6.1. "The **Riemann-Liouville integral** operator of order $\alpha > 0$ of a function $g(t), t \in (c, d)$ is represented as:

$$_{RL}\mathcal{J}^{\alpha}_{c,t}g(t) = \frac{1}{\Gamma(\alpha)} \int_{c}^{t} (t-u)^{\alpha-1}g(u)du, \qquad (6.1.2)$$

where $\Gamma(\cdot)$ is Euler's gamma function."

Definition 6.2. "The **Riemann-Liouville derivative** of order $\alpha > 0$ of a function $g(t), t \in (c, d)$ is defined as:

$${}_{RL}\mathcal{D}^{\alpha}_{c,t}g(t) = \frac{1}{\Gamma(\zeta - \alpha)} \frac{d^{\zeta}}{dt^{\zeta}} \int_0^t (t - u)^{\zeta - \alpha - 1} g(u) du, \qquad (6.1.3)$$

 $\zeta - 1 < \alpha < \zeta, \ \zeta \in \mathbb{N}.$ In particular, for $0 < \alpha < 1$, we have n = 1, and hence,

$${}_{RL}\mathcal{D}^{\alpha}_{c,t}g(t) = \frac{1}{\Gamma(1-\alpha)}\frac{d}{dt}\int_{c}^{t}(t-u)^{-\alpha}g(u)du.$$
 (6.1.4)

Definition 6.3. "The **Caputo fractional derivative** of order $\alpha > 0$ of a function $g(t), t \in (c, d)$ is defined as:

$${}_{C}\mathcal{D}^{\alpha}_{c,t}g(t) = \frac{1}{\Gamma(\zeta - \alpha)} \int_{c}^{t} (t - u)^{\zeta - \alpha - 1} g^{(\zeta)}(u) du, \qquad (6.1.5)$$

 $\zeta - 1 < \alpha < \zeta, \ \zeta \in \mathbb{N}.$ In particular, for $0 < \alpha < 1$, we have n = 1, and hence,

$${}_{C}\mathcal{D}^{\alpha}_{c,t}g(t) = \frac{1}{\Gamma(1-\alpha)} \int_{c}^{t} (t-u)^{-\alpha} g'(u) du.$$
 (6.1.6)

To apply the Haar wavelet following integral on Interval [0,1] is required :

$$\mathcal{P}_i^{\alpha}(t) = \frac{1}{\Gamma(\alpha+1)} \int_0^t (t-u)^{\alpha} \mathfrak{h}_i(u) du.$$

The R-L integration of (2.2.3) yields,

$$\mathcal{P}_{i}^{\alpha}(t) = \frac{1}{\Gamma(\alpha+1)} \begin{cases} \phi_{1}(t) & t \in [\vartheta_{1}(i), \vartheta_{2}(i)), \\ \phi_{2}(t) & t \in [\vartheta_{2}(i), \vartheta_{3}(i)), \\ \phi_{3}(t) & t \in [\vartheta_{3}(i), 1), \\ 0 & otherwise, \end{cases}$$
(6.1.7)

$$\begin{split} \phi_1(t) &= (t - \vartheta_1(i))^{\alpha}, \\ \phi_2(t) &= [(t - \vartheta_1(i))^{\alpha} - 2(t - \vartheta_2(i))^{\alpha}], \\ \phi_3(t) &= [(t - \vartheta_1(i))^{\alpha} - 2(t - \vartheta_2(i))^{\alpha} + (t - \vartheta_3(i))^{\alpha}], \\ \text{where } \vartheta_1(i) &= k/2^i, \vartheta_2(i) = (k + 0.5)/2^j, \vartheta_3(i) = (k + 1)/2^j. \end{split}$$

The index $i = 2^j + k + 1$, j = 0, 1, ..., J, where J is maximum level of wavelet and $k = 0, 1, ..., 2^j - 1$.

6.2 Description of Method

In this section we present Haar wavelet series method (HWSM) to find the approximate solution of the proportional delay Riccati differential equation of fractional order represented in equation (6.1.1). For that, we approximate $\chi^{\alpha}(t)$ present in equation (6.1.1) by truncated Haar wavelet series as follows:

$$\chi^{\alpha}(t) = \sum_{i=1}^{2^{J+1}} a_i \mathfrak{h}_i(t).$$
(6.2.1)

R-L Integration of (6.2.1) from 0 to t yields,

$$\chi(t) = \sum_{i=1}^{2^{J+1}} a_i \mathcal{P}_i^{\alpha}(t) + \chi(0).$$
(6.2.2)

Now replace t by qt in equation (6.2.2), we get,

$$\chi(qt) = \sum_{i=1}^{2^{J+1}} a_i \mathcal{P}_i^{\alpha}(qt) + \chi(0).$$
(6.2.3)

Using equations (6.2.1) to (6.2.3) in equation, (6.1.1) we get,

$$\sum_{i=1}^{2^{J+1}} a_i \mathfrak{h}_i(t) = \Omega(t) + c_1(t) (\sum_{i=1}^{2^{J+1}} a_i \mathcal{P}_i^{\alpha}(t) + \chi(0)) + c_2(t) (\sum_{i=1}^{2^{J+1}} a_i \mathcal{P}_i^{\alpha}(qt) + \chi(0)) (c_3(t) - (\sum_{i=1}^{2^{J+1}} a_i \mathcal{P}_i^{\alpha}(qt) + \chi(0))).$$
(6.2.4)

Discretize the system (6.2.4) with the chosen collocation points $t_l = \frac{(l-0.5)}{2^{J+1}}$, we get,

$$\sum_{i=1}^{2^{J+1}} a_i \mathfrak{h}_i(t_l) = \Omega(t_l) + c_1(t_l) (\sum_{i=1}^{2^{J+1}} a_i \mathcal{P}_i^{\alpha}(t_l) + \chi(0)) + c_2(t_l) (\sum_{i=1}^{2^{J+1}} a_i \mathcal{P}_i^{\alpha}(qt_l) + \chi(0)) (c_3(t_l) - (\sum_{i=1}^{2^{J+1}} a_i \mathcal{P}_i^{\alpha}(qt_l) + \chi(0))).$$
(6.2.5)

Solve the above system for Haar wavelet coefficients a_i 's. Plugging these coefficients into the equation (6.2.2) produces the approximate solution $\chi(t_l)$.

6.3 Applications and Numerical Results

The combination of fractional calculus with the theory of delay differential equations has enhanced the mathematical description of a number of real-world phenomena during the past few years. On the other hand, several numerical treatment has been developed for solving fractional differential models. However, very few researchers have thoroughly investigated fractional differential equations with delay.

In this section, we shall be concerned with numerical treatment of some fractional order delay differential equations using Haar wavelet series method (HWSM). Before solving numerical examples, we shall state some real world applications of HWSM from existing literature. In Dec-2019 a threatful outbreak called the novel corona virus-2019 disease brought the world to its knees and took daily life to a grinding halt in much of the world. The researchers claim that the virus was initiated in the Chinese city of Wuhan. Planet-wide research to identify the symptoms, to control its spread, and to cure & eradicate the disease is still in full swing. In an attempt kamal shah et al. [98] study the transmission dynamics of the novel coronavirus-2019 and construct a fractional order differential mathematical model by considering three compartments including the susceptible population, infected population, and recovered population. Further, the solution of the model is computed using the Haar wavelet collocation method. Hence the method is proven as efficient tool in infectious disease spread modeling. Several recent studies which have promoted Haar wavelet as favorable mathematical tool are [57, 73, 102].

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FIGURE 6.1: Problem 1.

Problem 1: Consider the following fractional order PDRDE

$$y^{\alpha}(t) = \frac{1}{4}y(t) + y(\frac{t}{2})(1 - y(\frac{t}{2})), 0 \le \alpha \le 1, t \in (0, 1),$$
(6.3.1)

with initial condition y(0) = 1, and possesses the exact solution

$$y(t) = \frac{1}{2} + \frac{1}{2}\cos(\frac{\sqrt{2}t}{4}) + \frac{\sqrt{2}}{2}\sin(\frac{\sqrt{2}t}{4})$$

when $\alpha = 1$.

The approximate solution of (6.3.1) is computed using HWSM. The solution behaviour at $\alpha = 0.5, 0.6, 0.7, 0.8, 0.9, 1$ is presented graphically in Fig. 6.1. Also, we have presented solutions for $\alpha = 0.5, 0.7, 0.9, 1$ at selected collocation points in Table 6.1. Maximum absolute errors (MAEs) at different wavelet levels J are demonstrated in Table 6.4.

Problem 2: Solve the following fractional order PDRDE

$$y^{\alpha}(t) = \frac{1}{2}exp(\frac{t}{2})y(\frac{t}{2}) + \frac{1}{2}y(t), 0 \le \alpha \le 1, t \in (0,1)$$
(6.3.2)
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t_l	$\alpha = 0.5$	$\alpha = 0.7$	$\alpha = 0.9$	$\alpha = 1$	y_{exact}
0.0313	1.0455	1.0235	1.0113	1.0078	1.0078
0.1563	1.0933	1.0692	1.0473	1.0383	1.0383
0.2813	1.1173	1.1000	1.0784	1.0677	1.0677
0.4063	1.1335	1.1243	1.1065	1.0960	1.0961
0.5313	1.1456	1.1442	1.1322	1.1232	1.1232
0.6563	1.1551	1.1609	1.1558	1.1492	1.1492
0.7813	1.1627	1.1751	1.1775	1.1739	1.1739
0.9063	1.1689	1.1872	1.1973	1.1972	1.1973

TABLE 6.1: Approximate solution at $\alpha = 0.5, 0.7, 0.9, 1$ when J=3 (Problem 1).



FIGURE 6.2: Problem 2.

with initial condition y(0) = 1, and has the exact solution y(t) = exp(t) when $\alpha = 1$.

The approximate solution of (6.3.2) is computed using HWSM. The solutions behavior at $\alpha = 0.5, 0.6, 0.7, 0.8, 0.9, 1$ is presented graphically in Fig. 6.2. Also, we have presented solutions for $\alpha = 0.5, 0.7, 0.9, 1$ at selected collocation points in Table 6.2. Maximum absolute errors (MAEs) at different wavelet levels J are demonstrated in table 6.4.

Problem 3: Solve the following fractional order PDRDE

$$y^{\alpha}(t) = -\frac{1}{8}y(t) + y(\frac{t}{2})(1 - y(\frac{t}{2})), 0 \le \alpha \le 1, t \in (0, 1),$$
(6.3.3)

t_l	$\alpha = 0.5$	$\alpha = 0.7$	$\alpha = 0.9$	$\alpha = 1$	y_{exact}	
0.0313	1.2426	1.1065	1.0480	1.0323	1.0317	
0.1563	1.6442	1.3661	1.2183	1.1697	1.1691	
0.2813	1.9987	1.6148	1.3998	1.3255	1.3248	
0.4063	2.3653	1.8796	1.6007	1.5021	1.5012	
0.5313	2.7606	2.1698	1.8255	1.7021	1.7011	
0.6563	3.1951	2.4919	2.0782	1.9288	1.9276	
0.7813	3.6775	2.8520	2.3632	2.1857	2.1842	
0.9063	4.2165	3.2561	2.6849	2.4768	2.4750	

TABLE 6.2: Approximate solution at $\alpha = 0.5, 0.7, 0.9, 1$ when J=3 (Problem 2).

TABLE 6.3: Approximate solution at $\alpha = 0.5, 0.7, 0.9, 1$ when J=3 (Problem 3).

t_l	$\alpha = 0.5$	$\alpha = 0.7$	$\alpha = 0.9$	$\alpha = 1$	y_{exact}	
0.0313	0.2825	0.2655	0.2572	0.2549	0.2549	
0.1563	0.3248	0.2985	0.2810	0.2747	0.2746	
0.2813	0.3522	0.3243	0.3031	0.2947	0.2947	
0.4063	0.3743	0.3472	0.3246	0.3150	0.3150	
0.5313	0.3935	0.3684	0.3457	0.3355	0.3355	
0.6563	0.4106	0.3883	0.3665	0.3562	0.3562	
0.7813	0.4263	0.4073	0.3871	0.3770	0.3770	
0.9063	0.4407	0.4255	0.4075	0.3981	0.3981	

with initial condition $y(0) = \frac{1}{4}$, and the exact solution

$$y(t) = \frac{1}{2} - \frac{1}{4} cos(\frac{\sqrt{5}t}{8}) + \frac{\sqrt{5}}{4} sin(\frac{\sqrt{5}t}{8})$$

when $\alpha = 1$.

The approximate solution of (6.3.3) is computed using HWSM. The Solution behavior at $\alpha = 0.5, 0.6, 0.7, 0.8, 0.9, 1$ is presented graphically in Fig. 6.3. Also, we have presented solution for $\alpha = 0.5, 0.7, 0.9, 1$ at selected collocation points in table 6.3. Maximum absolute errors (MAEs) at different wavelet levels J are demonstrated in Table 6.4.

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FIGURE 6.3: Problem 3.

Гавle 6.4: Max	imum absolute	$\operatorname{error}(\operatorname{MAE})$	$max y_{approx}$	$-y_{exact}$
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J	Problem 1	Problem 2	Problem 3	Problem 4
3	3.0351E - 05	1.8900E - 03	9.4852E - 06	2.8315e - 08
4	7.6090E - 06	4.8189E - 04	2.3778E - 06	1.1102e - 16
5	1.9048E - 06	1.2167E - 04	5.9526E - 07	1.1102e - 16
6	4.7652E - 07	3.0572E - 05	1.4891E - 07	—
7	1.1917E - 07	7.6623E - 06	3.7240E - 08	—
8	2.9797E - 08	1.9178E - 06	9.3117E - 09	_
9	7.4499E - 09	4.7979E - 07	2.3281E - 09	—
10	1.8620E - 09	1.1996E - 07	5.8205E - 10	—

Problem 4: Now for comparison we choose the following fractional Riccati differential equation from literature [94]

$$y^{\alpha}(t) = t^{3}y^{2}(t) - 2t^{4}y(t) + t^{5}, 0 \le \alpha \le 1, t \in (0, 1)$$
(6.3.4)

with initial condition y(0) = 0, and has the exact solution y(t) = t when $\alpha = 1$. The aforementioned problem has been solved by means of HWSM and FDE solver. It is noted that the HWSM provides better accuracy and converges towards an exact solution, which has been asserted with the help of the graphical plot 6.4. To analyse the solution behaviour, a tabular (6.5) comparison is performed with $\alpha = 0.5, 0.6, 0.7, 0.8, 0.9, 1$. Also, maximum absolute errors and a comparison with IRKHSM are presented in table 6.4 and table 6.7.



FIGURE 6.4: Problem 4.

TABLE 6.5: Approximate solution at $\alpha = 0.5, 0.7, 0.9, 1$ (Problem 4)).
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t_l	$\alpha = 0.5$	$\alpha = 0.7$	$\alpha = 0.9$	$\alpha = 1$	y_{exact}
0.0313	0.1995	0.0973	0.0460	0.0313	0.0313
0.1563	0.4461	0.3001	0.1956	0.1563	0.1563
0.2813	0.5990	0.4530	0.3320	0.2813	0.2813
0.4063	0.7215	0.5863	0.4622	0.4063	0.4063
0.5313	0.8277	0.7080	0.5885	0.5313	0.5313
0.6563	0.9236	0.8220	0.7119	0.6562	0.6563
0.7813	1.0115	0.9300	0.8329	0.7812	0.7813
0.9063	1.0920	1.0332	0.9521	0.9062	0.9063

TABLE 6.6: Comparison of MAE(Problem 4).

HWSM	HWSM	IRKHSM
J = 3	J = 4	N = 4[94]
2.8315E - 08	1.1102E - 16	2E - 07

	$\alpha = 1$	MAE	fde	Solver[37]	0	2.011643E - 07	1.132459E - 06	1.132459E - 06	1.132459E - 06	1.132459E - 06	1.132460E - 06	1.132464E - 06	1.132477E - 06	1.132511E - 06	1.132594E - 06	1.132784E - 06	1.133185E - 06	1.133985E - 06	1.135503E - 06	1.138263E - 06
	$\alpha = 1$	MAE	HWSM		0	0	0	0	0	0	0	0	0	1.998401E - 15	4.596323E - 14	1.536992E - 12	2.236499 E - 11	2.876829E - 10	3.129013E - 09	2.831572E - 08
lver .	$\alpha = 1$	fde	Solver		0.031250	0.093749	0.156248	0.218748	0.281248	0.343748	0.406248	0.468748	0.531248	0.593748	0.656248	0.718748	0.781248	0.843748	0.906248	0.968748
vs FDE So	$\alpha = 1$	MSWH			0.031250	0.093750	0.156250	0.218750	0.281250	0.343750	0.406250	0.468750	0.531250	0.593749	0.656249	0.718749	0.781249	0.843749	0.906249	0.968749
n of HWSM	$\alpha =$	0.9	fde	Solver	0.031250	0.116997	0.191261	0.261731	0.329847	0.396267	0.461370	0.525402	0.588534	0.650894	0.712579	0.773664	0.834211	0.894268	0.953874	1.013060
Comparisor	$\alpha =$	0.9	MSWH		0.045951	0.123510	0.195600	0.264780	0.331986	0.397705	0.462240	0.525794	0.588514	0.650511	0.711872	0.772664	0.832943	0.892753	0.952131	1.011104
ABLE 6.7:	$\alpha =$	0.7	fde	Solver	0.031250	0.189275	0.287972	0.372261	0.448408	0.519067	0.585689	0.649179	0.710147	0.769025	0.826125	0.881675	0.935838	0.988726	1.040415	1.090949
E	$\alpha =$	0.7	MSWH		0.097275	0.209889	0.300120	0.379851	0.452972	0.521399	0.586262	0.648303	0.708035	0.765831	0.821966	0.876642	0.930006	0.982160	1.033172	1.083086
	$\alpha =$	0.5	fde	Solver	0.031250	0.313356	0.430277	0.520157	0.596201	0.663572	0.724935	0.781869	0.835373	0.886100	0.934470	0.980737	1.025027	1.067373	1.107743	1.146071
	$\alpha =$	0.5	HWSM		0.199471	0.345504	0.446102	0.528002	0.599044	0.662854	0.721469	0.776156	0.827741	0.876778	0.923624	0.968496	1.011508	1.052690	1.092022	1.129461
	t_l				0.03125	0.09375	0.15625	0.21875	0.28125	0.34375	0.40625	0.46875	0.53125	0.59375	0.65625	0.71875	0.78125	0.84375	0.90625	0.96875

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FIGURE 6.5: Absolute error at wavelet level J = 4 (Problem 4).

6.4 Summary

In this chapter¹, a collocation method based on the delayed Haar basis and their fractional integrals is discussed for the approximation of fractional proportional delay Riccati differential equations. A short account of fractional calculus and its applications was also presented. The applicability and utility of the method are tested by solving a few benchmark problems. The results we obtained are very promising.

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Concluding Remarks

The work presented in this thesis has focused on establishing accurate, stable, and efficient numerical techniques for solving a particular class of delay differential equations known as proportional delay differential equations.

A numerical technique, presented in chapter 2, was designed to solve proportional delay Riccati differential equations. In this technique, Haar wavelets are preferred primarily due to their low computational cost and mathematical simplicity. Several examples have been solved to demonstrate the efficiency and accuracy of the technique. Numerical comparisons have been made between the implementations of the proposed method, the Bezier Curves Method, and the exact solution. The numerical experiment indicates that the Haar wavelet method has given an accuracy which varies between $O(10^{-05}) - O(10^{-09})$ for the different resolutions of the wavelet. The method is extended and developed to solve Dirichlet BVPs of proportional delay nature and simultaneous proportional delay differential equations, as presented in chapters 3 and 4, respectively. The method has proven to be an efficient mathematical tool for solving BVPs because it can automatically incorporate boundary conditions. Several benchmark problems have been solved, maximum absolute errors and convergence rates are calculated. The accuracy for BVPs varies between $O(10^{-04}) - O(10^{-10})$ when J varies from 1 to 10 (Table 3.1). Also, the convergence rate for proportional delay IVPs and BVPs matches closely with the theoretical value, i.e., 2 (Table 3.2). Similarly, the method is extended to solve simultaneous proportional delay differential equations. Some linear and non-linear problems with two and three unknown functions are solved, and a comparison has been made between the solution obtained via the Haar wavelet series method and the exact solution. Figure 4.3 and 4.5 show that an increase in resolution

contributes to the decrease in maximum absolute errors, and the method performs well even for the system of equations.

In chapter 5, a modified Haar wavelet approach is established to investigate the solution of higher-order multipanotgraph DEs arising in electrodynamics. In this modified approach, instead of the highest derivative (n^{th}) functions present in differential equations, the $(n + 1)^{th}$ derivative of the unknown function is expanded into the Haar wavelets series with unknown coefficients, and the other derivatives are obtained through integration. This approach transforms the problem into a set of algebraic equations and thus simplifies the solution process. Numerical results show that the method is effective, robust, and comparable with existing methods (see, Table 5.1, 5.7), such as the Two-stage RKM, One-leg θ , VIM, and RKHS method.

Finally, in chapter 6, fractional integrals of Haar wavelets in the R-L sense are defined, and then an efficient Haar wavelets-based algorithm is established to obtain the numerical solution of fractional order proportional delay Riccati DEs. The outcomes of the numerical simulations in Table 6.6 and 6.7 indicate the superiority of HWSM over IRKHSM and FDESOLVER. The accuracy varies between $O(10^{-05}) - O(10^{-10})$ (Table 6.4) when J varies from 3 to 10, and in one case, it varies between $O(10^{-08}) - O(10^{-16})$ (Table 6.7).

A possible future research direction could be:

A Numerical method for solving delay fractional order partial DEs.

To design and investigate fractional delay differential models using wavelets.

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